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Informal Report

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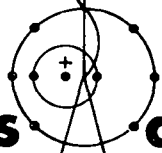
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**Methods for Calculating Group Cross Sections for
Doubly Heterogeneous Thermal Reactor Systems**



Issued: February 1977



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Methods for Calculating Group Cross Sections for Doubly Heterogeneous Thermal Reactor Systems

by

M. G. Stamatelatos
R. J. LaBauve



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METHODS FOR CALCULATING GROUP CROSS SECTIONS
FOR DOUBLY HETEROGENEOUS THERMAL REACTOR SYSTEMS

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M. G. Stamatelatos and R. J. LaBauve

ABSTRACT

This report discusses methods used at LASL for calculating group cross sections for doubly heterogeneous HTGR systems of the General Atomic design. These cross sections have been used for the neutronic safety analysis calculations of such HTGR systems at various points in reactor lifetime (e.g., beginning-of-life, end-of-equilibrium cycle). They were also compared with supplied General Atomic cross sections generated with General Atomic codes. The overall agreement between the LASL and the GA cross sections has been satisfactory.

I. INTRODUCTION

Over approximately the past two and one-half years, the Los Alamos Scientific Laboratory has been engaged in reactor safety studies for High Temperature Gas-cooled Reactor (HTGR) systems of the General Atomic design. Discussed in this report is the methodology connected with a small part of this effort, namely the calculation of multigroup cross sections for use in neutronic calculations (e.g., effective multiplication factors, temperature coefficients, etc.). The initial effort has been directed towards using generally available computer codes with minimal effort in the direction of new methods development. Unfortunately, however, many specialized GA codes were kept proprietary and other widely available codes were not specialized enough to correctly treat special configurations like, for example, doubly heterogeneous HTGR systems. Therefore, at some point in the cross-section development, it was decided to intensify the development of methods to treat such system peculiarities. Therefore, as it

will be seen in the following discussion, the final code system configuration used resembles little the initial configuration used for calculating homogeneous HTGR cross sections.

II. HOMOGENEOUS CROSS SECTIONS

In the initial stages of the cross-section generation process, a number of code systems were explored and these are discussed here mostly for the sake of "historic" completeness. Although these systems are quite different from the final system used, they are nevertheless valid options for generating homogenized-medium cross sections or cross sections for media with one allowed level of heterogeneity. Approximate ways of incorporating the effects of the second level of heterogeneity (fuel grains in a fuel rod) have been explored, as will be seen later, but the final system chosen has proved to be superior to the others in all respects including accuracy and flexibility.

The initial data flow system (including options) for generating homogeneous-medium few-group cross sections is shown in Fig. 1. The starting point has always been the basic Evaluated Nuclear Data Files (ENDF/B) cross sections (initially version III; later several version IV elements were included). The few-group neutron energy structure used in all the work described in this report has been a nine-group General Atomic structure (adopted for comparison with supplied GA cross sections) shown in Table I. The initial set of tempera-

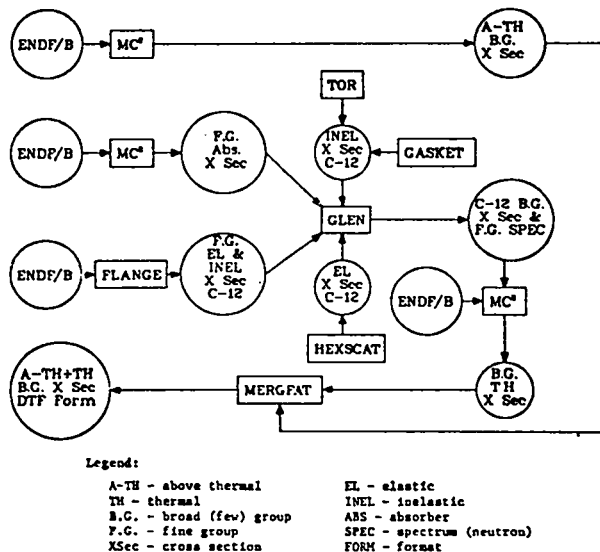


Fig. 1.
Initial data flow systems (several options are shown).

TABLE I

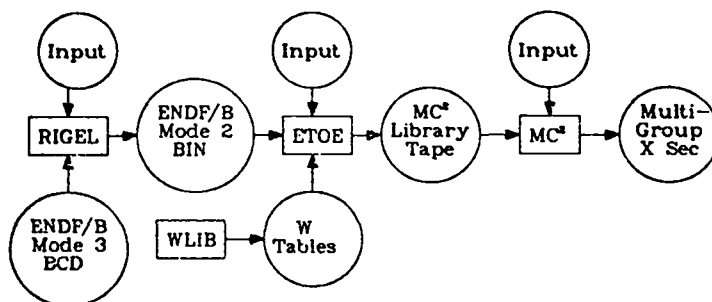
FEW-GROUP ENERGY STRUCTURE
 $E_{\max} = 10 \text{ MeV}$

Group No.	Lower Boundary (eV)
1	$1.83 \times 10^{+5}$
2	$9.61 \times 10^{+2}$
3	$1.76 \times 10^{+1}$
4	3.93
5	2.38
6	4.14×10^{-1}
7	1.00×10^{-1}
8	4.00×10^{-2}
9	5.00×10^{-4}

tures for which few-group cross sections were generated is: 300, 500, 800, 1200, 1700, 2300, and 3000 K. These were used for the beginning-of-life (BOL) composition. Later, several other temperatures (600, 1000, 1500, 2000, and 2600 K) were also included for a more accurate evaluation of the temperature coefficient at the end-of-equilibrium-cycle (EOEC) composition.

The above-thermal (10 MeV - 2.38 eV) cross sections of the system shown in Fig. 1 were generated with an operational LASL-modified version of MC²-I¹ code that requires special library preparation, i.e., it does not directly operate on the ENDF/B cross-section files. The preparation of such an MC² input file is shown in the diagram of Fig. 2. The RIGEL² code is used to convert ENDF/B data in standard BCD format (Mode 3) to an alternate binary format (Mode 2). The ETOE³ code prepares a library tape for MC² including "W-tables" that are supplied by the WLIB code. Since ETOE provides pointwise elastic-scattering cross sections for MC², temperature must be an input parameter to ETOE which means that a different MC² library tape must be prepared for each temperature. The various MC² libraries are then merged with an auxiliary code, MERMC2, not shown in Fig. 1. There are certain limitations connected with the MC² code, some of which have proved to be so hard to circumvent, unless considerable effort was put in modifying the code, that MC²-I was removed from the final data flow system to be discussed later. First, because of storage limitations, fine-group cross sections for the entire energy range (10 MeV - 10⁻⁵ eV) cannot be generated in one pass, so that separate but slightly overlapping problems were run for the "high" (10 MeV - 0.414 eV) and "low" (2.38 - 5 x 10⁻⁴ eV) energy ranges. Second, the maximum energy value in MC²-I is fixed (10 MeV) and one is also forced to use a fixed-lethargy grid in one of two available options, "all-fine" with Δu = 0.25 and "ultra-fine" with Δu = 1/120. Since the second option was found to be too time-consuming and costly without the benefit of considerable increase in

Fig. 2.
MC²-I library preparation.



output cross-section quality, the "all-fine" option was chosen for generating both above-thermal and thermal fine-group cross sections in the GAM-I constant-lethargy structure of 0.25. The spectrum-weighting function specified for the derivation of fine-group cross sections was chosen to be $1/E$ for the above-thermal region and a "properly hardened Maxwellian" for the thermal region. The latter was calculated by the thermal code GLEN.⁴

The graphite cross sections in the thermal region were treated separately. Initially, the FLANGE⁵ code was used to interpolate (both energy-wise and temperature-wise) preprocessed graphite thermal inelastic-scattering cross sections available in ENDF/B format (MAT 1065, MF 4 and 7). This process has proved costly and inefficient by comparison with directly calculating the $S(\alpha, \beta)$ data from codes like GASKET⁶ or TOR.⁷ The graphite coherent elastic cross section was calculated with a modified version of the HEXSCAT⁸ code which now calculates Legendre elastic-scattering components up to the order 5. All fine-group thermal cross sections were collapsed with the GLEN code to the required few-group set. MC² was used to collapse the above-thermal fine-group cross sections to the corresponding few-group set. An auxiliary code MERGFAT (Appendix C) was used to merge the fast and thermal few-group cross sections in the proper format required by the DTF-IV⁹ neutronics transport code.

Several modifications to MC²-I were made. An important one was in the multigroup averaging method for the resolved-resonance capture cross sections. The MC²-I method is given by the following equations:

$$\left(\frac{\sigma_{cJ}}{\sigma}\right)_{\text{resolved}}^{\text{bg}} = \frac{\sum_{j \text{ in } J} \langle \sigma \rangle_j^{\text{fg resolved}} Q_j}{\sum_{j \text{ in } J} Q_j} \quad , \quad (1)$$

where

$$Q_j = \int_{E_j}^{E_{j+1}} S_j^{\text{fg}} \frac{dE}{\Sigma_T(E)} \quad , \quad (2)$$

and

$$\begin{aligned}
S_j^{fg} = & \sum_{k \neq j} \Sigma_{inel,k}^{fg} \phi_k^{fg} \frac{\langle \Sigma_{inel} \rangle_{k \rightarrow j}}{\langle \Sigma_{inel} \rangle_k} + \sum_{k \neq j} \Sigma_{n,2n,k}^{fg} \phi_k^{fg} \frac{2 \langle \Sigma_{n,2n} \rangle_{k \rightarrow j}}{\langle \Sigma_{n,2n} \rangle_k} \\
& + \langle \Sigma_{el} \rangle_{j-1 \rightarrow j} \phi_{j-1}^{fg} \quad ,
\end{aligned}
\tag{3}$$

where superscripts fg and bg indicate fine-group and broad-group, respectively. J and j are subscripts referring to broad-group and fine-group, respectively. This method of averaging has produced unsatisfactory results and, since it had not been shown to be valid for thermal reactor systems, it was replaced by the usual spectrum-weighting method used by almost all multigroup codes. This change has resulted in much better MC²-I cross sections.

Regarding other codes used, one of the most important changes was made in GLEN whose original version did not allow for energy-dependent scattering cross sections of nonmoderator materials. Although for most heavy absorbers it is possible to give the thermal-scattering cross sections in terms of an average energy-independent number, some resonance elements like ¹³⁵Xe or ¹⁴⁹Sm definitely require energy-dependent scattering cross sections. Therefore, modifications were made in the GLEN code to allow the option of including energy-dependent scattering cross sections together with the only previously available option of supplying a single energy-independent scattering cross section value for each nonmoderator material. The choice of options in the modified GLEN version is made by means of a flag, ISCAT.

The homogeneous cross sections produced by the scheme of Fig. 1, including all discussed modifications, for a beginning-of-life HTGR composition were found good* and the discrepancies between these cross sections and the supplied General Atomic (GA) cross sections produced with the proprietary MICROX¹⁰ code were in the direction attributable to heterogeneity effects or to different initial basic data. The incorporation of double-heterogeneity effects by the MICROX method (for comparison with the GA cross sections) was found to essentially amount to reprogramming the MC²-I code. Although alternate approximate methods of incorporating double heterogeneity effects in codes like MC²-I were developed, as

*Comparisons were made with the MC²-II code, courtesy of H. Henryson of ANL.

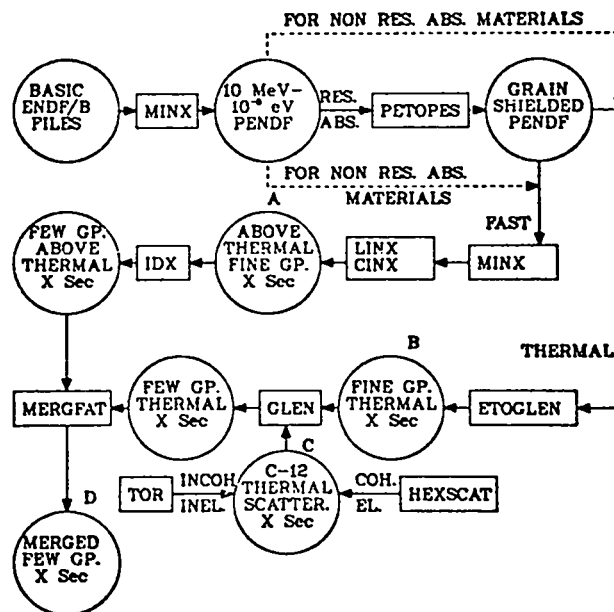


Fig. 3.
Final data flow system.

it will be discussed later, we have decided to adopt a totally new data flow system (Fig. 3), more modern and more flexible including the MINX¹¹ code which was developed at LASL.

III. CROSS SECTIONS FOR DOUBLY HETEROGENEOUS HTGR SYSTEMS

The latest version of the data flow system (Fig. 3) also starts from the basic ENDF/B file. The MINX code generates temperature-broadened pointwise cross sections in the ENDF/B format (PENDF) and further collapses them to the desired fine-group structure in the Bondarenko¹² energy-shielding formalism. For resonance absorber materials, the PENDF cross sections are space shielded over the entire energy range according to the Wälti formalism¹³ adopted in the GA code MICROX to account for the grain heterogeneity in HTGR fuel rods. For this purpose, a special code, PETOPES (Appendix A), was written. The fast-group cross sections were then collapsed by the MINX code to a 69-group fine-group structure (68 equal-lethargy groups from 10 MeV to 0.414 eV, plus 1 dump group) and further collapsed by the IDX¹⁴ code to the desired broad-group structure (see Table I). Corrections for the second level of heterogeneity (fuel rods in the reactor core) were applied by the rational-approximation collision-probabil-

ity method of Levine¹⁵ in a modified version of the LDX code that can handle the Bondarenko formalism provided by the MINX code.

The thermal portion of the PENDF cross sections was processed by a specially written code, ETOGLEN (Appendix B), and by the GLEN code. Since GLEN requires pointwise rather than groupwise cross sections, ETOGLEN was written to select a thermal fine-group structure in such a way as to best calculate resonance integrals by the GLEN method. GLEN calculates a properly hardened thermal neutron spectrum based on the input isotopic composition and collapses the fine-group (points) cross sections to the required few-group thermal structure. GLEN also accepts graphite elastic-scattering cross sections and scattering-law data as calculated, for example, by the HEXSCAT and TOR codes, respectively. MERGFAT (Appendix C) was used to merge the fast and thermal few-group cross sections in the required DTF-IV format.

The operation of the code system shown in Fig. 3 proceeds as follows:

1. Using the basic ENDF/B file as input, a pointwise ENDF/B file (PENDF) is prepared by the MINX code for each nuclide needed in the neutronic calculations. Nuclides prepared for the HTGR composition are shown in Table II. The data in the PENDF files are given at 0, 300, 950, and 3000 K.
2. If the cross sections of a nuclide are not to be grain shielded, the PENDF file is processed directly by the MINX code to give 69-group cross sections for input to the LINX-CINX codes.^{16,17} The 69-group structure consists of the GAM-I group structure plus a dump group necessary to obtain the correct eigenvalue in LDX. The weighting function used in MINX for generating the 69-group set is shown in Fig. 4. It is the composite result of calculations for a typical HTGR system made with the GLEN and MC² codes.
3. For those nuclides for which grain shielding is important, the PENDF files are used as input to the PETOPES code, which generates a grain-shielded PENDF file (PENDFS). This file is then used as input to MINX to generate multigroup cross sections as indicated in 2 above.
4. The LINX-CINX codes are used to combine multigroup data for all nuclides into a single data library used for input to the LDX code. This is the file designated by "A" in Fig. 3. Note that file A contains temperature-dependent f-factors for Bondarenko treatment by LDX.
5. The ETOGLEN code is used to retrieve cross-section thermal data (from 5×10^{-4} to 2.38 eV in the group structure of Table I) from the PENDF or PENDFS file for each nuclide and to supply pointwise cross sections for elastic scatter-

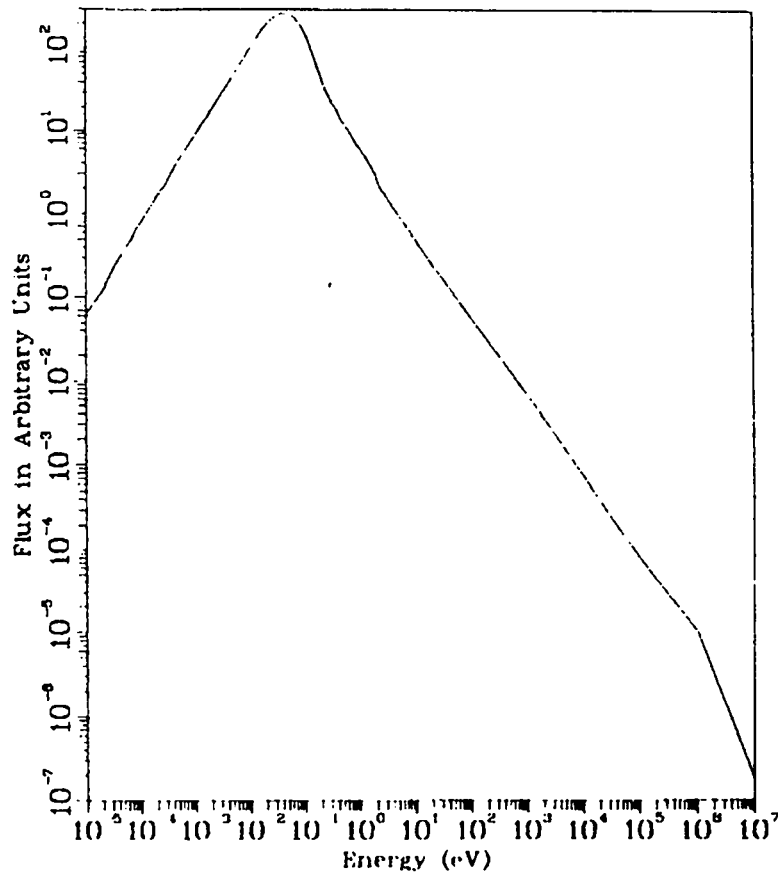


Fig. 4.
 Typical HTGR spectrum used for MINX weight function.

ing, nu times fission, and absorption cross sections for the GLEN code. Data for all nuclides at several temperatures (300, 950, 3000 K for HTGR) are combined to form data file "B" in Fig. 3.

6. Data file "C" (graphite in the case of an HTGR) is made by combining the outputs of the TOR and HEXSCAT codes into a single file. For graphite, crystal-lattice parameters are input to HEXSCAT and a phonon distribution (Young-Koppel) is input to TOR. Library "C" contains data for each temperature of interest (see Table II for an HTGR).
7. The final broad-group cross sections for all nuclides at a single temperature, file "D" in Fig. 3, are created using the code MERGFAT to merge the outputs of GLEN and LDX. This is usually done in a single run for

TABLE II

	<u>Nuclide</u>	<u>MAT NO.</u>	<u>ENDF/B- VERSION</u>	<u>Region</u>
1.	B-10	1155	III	Core
2.	C-12	1165	III	"
3.	O-16	1134	III	"
4.	Si-28	1194	III	"
5.	Xe-135	1294	IV	"
6.	Sm-149	1027	III	"
7.	Th-232	1117	III	"
8.	Pa-233	1119	III	"
9.	Pa-233	1297	IV	"
10.	U-233	1260	IV	"
11.	U-234	1043	III	"
12.	U-235	1157	III	"
13.	U-236	1163	III	"
14.	U-238	1158	III	"
15.	Pu-238	1050	III	"
16.	Pu-239	1264	IV	"
17.	Pu-240	1265	IV	"
18.	Pu-241	1266	IV	"
19.	Pu-242	1161	III	"
20.	B-10	1155	III	reflector
21.	C-12	1165	III	reflector

Cross sections for every nuclide in the above list are available for 12 temperatures: 300, 500, 600, 800, 1000, 1200, 1500, 1700, 2000, 2300, 2600, and 3000 Kelvin.

efficiency purposes. As cross sections at additional temperatures are generated, the data are added to the broad-group cross-section library by means of the UPDATE feature of the LASL CDC-7600 operating software.

IV. DOUBLE-HETEROGENEITY SPACE SHIELDING

Two methods of space shielding cross sections for a doubly heterogeneous reactor system are discussed here. The first method consists of the application of Wälti's¹³ method of grain shielding to pointwise (PENDF) cross sections followed by the application of the Levine¹⁵ formalism of "gross" (fuel-rod) space shielding to collapsed grain-shielded fine-group cross sections. The

grain shielding was implemented in the PETOPES code and the gross heterogeneity correction was made in a modified LDX code.

The second method of space shielding cross sections is a newly developed method based on rational approximations and collision probabilities which accounts for both levels of heterogeneity at the fine-group cross-section level. It, therefore, bypasses the time-consuming pointwise grain-shielding process and it serves as independent reference, since it produces results in close agreement with the first method.

A. First Method

1. Grain-Shielding Treatment. Wälti's grain-shielding method has been incorporated in the GA code MICROX and produces, according to Wälti's claims, results in close agreement with the detailed Nordheim integral method (NIT) used in the GAROL¹⁸ and the GGC-5¹⁹ codes.

In the Wälti procedure, the grain-shielded absorption cross section is given by

$$\sigma_i^{\text{eff}}(E) = \sigma_i(E) \frac{\Gamma(E)}{1 - r^3 [1 - \Gamma(E)]} \quad , \quad (4)$$

where

$\sigma_i(E)$ = unshielded energy-dependent cross section for the i -th heavy nuclide;

r = ratio of fuel-to-moderator radii in a two-concentric-sphere model (inner = fuel; outer = moderator) representing a uniform grain distribution in the fuel rod; and

$\Gamma(E)$ = self-shielding factor, i.e., the ratio of average neutron fluxes in the grain and in the moderator, $\bar{\phi}_0/\bar{\phi}_1$, where subscripts 0 and 1 refer to the grain and the surrounding moderator regions, respectively.

If, due to the presence of large amounts of moderator material, isotropic angular fluxes are assured for regions 0 and 1, the neutron balance equations for the two regions yield

$$\Gamma(E) = \frac{\bar{\phi}_0}{\bar{\phi}_1} = \frac{1 + \rho_Q [1 + W \bar{\ell}_1 (\Sigma_{a,1} + \Sigma_{\text{out},1})]}{1 + \rho_Q + W \bar{\ell}_0 (\Sigma_{a,0} + \Sigma_{\text{out},0})} \quad , \quad (5)$$

where

$$\rho = \frac{\bar{\lambda}_0}{\bar{\lambda}_1} = \frac{V_0}{V_1} = \text{volume ratio of regions 0 and 1,}$$

Q = ratio of spatially averaged source densities in regions 0 and 1,

$$W = 1 + \bar{H}_0(\Sigma_{t,0}) + \bar{H}_1(\Sigma_{t,1}) \quad , \quad (6)$$

$\bar{\lambda}_0, \bar{\lambda}_1$ = mean chord lengths in regions 0 and 1, respectively;

$$\bar{\lambda}_j = \frac{4V_j}{S_j} \quad , \quad j = 0,1 \quad .$$

The first-collision "augment" for region j, \bar{H}_j is given by

$$\bar{H}_j(\Sigma_{t,j}) = \frac{1 - \bar{P}_j}{\bar{\lambda}_j \bar{P}_j \Sigma_{t,j}} \quad , \quad j = 0,1 \quad , \quad (7)$$

and $\Sigma_{a,j}$, $\Sigma_{out,j}$, and $\Sigma_{t,j}$ are the macroscopic absorption, outscatter, and total group cross sections, respectively, for region j (0 or 1).

Augment $\bar{H}_1(\Sigma_t)$ can be approximated by $\bar{H}_1(0)$ which is given by the following expression

$$\begin{aligned} \bar{H}_1(0) = \left(\frac{\gamma}{r}\right)^2 & \left\{ (1-r^2)^2 \left(1 + \frac{1}{4} \ln \frac{1+r}{1-r}\right) - \frac{r}{2} (1-r)^2 \right. \\ & \left. + \left(\frac{2}{3r}\right)^2 \left[(1-r^2)^3 - 3(1-r^3)^2 + 2(1-r^3)(1-r^2)^{3/2} \right] \right\} \quad , \end{aligned} \quad (8)$$

where

$$\gamma = \frac{3r^2}{4(1-r^3)} \quad . \quad (9)$$

The escape probability function \bar{P}_0 is given by the expression of Case et al.²⁰

$$P_0(\Sigma_{t,0}) = \frac{3}{8X^3} [2X^2 - 1 + (1 + 2X) \exp(-2X)] \quad , \quad (10)$$

where

$$x = \frac{3}{4} \bar{\ell}_0 \Sigma_{t,0} \quad . \quad (11)$$

Source density ratio Q can be calculated from

$$Q = \frac{\xi_{0,pot} \Sigma_{S,0}^{pot}}{\xi_{1,pot} \Sigma_{S,1}^{pot}} \quad , \quad (12)$$

and the self-scattering cross section at the pointwise level is approximated by

$$\Sigma_{SS,j}(E) \approx \frac{1 - \xi_j(E)}{\xi_1^{pot}} \Sigma_{S,j}(E) \quad , \quad j = 0, 1 \quad , \quad (13)$$

where the average logarithmic energy decrement $\xi_j(E)$ is given by

$$\xi_j(E) = \frac{\sum_i \xi_j^i \Sigma_{S,j}^i(E)}{\Sigma_{S,j}(E)} \quad , \quad j = 0, 1 \quad , \quad (14)$$

i being the nuclide index.

The derivations of these equations and the justifications for the approximations made can be found in Wälti's paper.¹³ The above summary of the theory has been included only for readers' convenience. The programming of the equations in the PETOPES code is discussed in Appendix A.

2. Fuel-Rod Heterogeneity Treatment. The escape probability from a regular array of fuel (absorber) lumps, each assumed to be homogeneous in composition, is given by the Nordheim expression

$$P_{esc}^* = P_{esc} \frac{1 - C}{1 - C(1 - \Sigma_F \bar{\ell}_F P_{esc})} \quad , \quad (15)$$

where

P_{esc} = escape probability from one lump,
 C = Dancoff factor (Appendix D), and
 \bar{l}_F = fuel-rod mean chord length.

Equations for P_{esc} for different lump geometries have been derived by many investigators (e.g., see Refs. 20, 21, 22). Wigner²³ has proposed a "rational" approximation to P_{esc} which gives the correct value in the two limiting cases of very large and very small lumps. For better approximations between these two extreme limits, various Wigner-like approximations have been proposed. One such popular approximation is due to Levine¹⁵ and is given by the following expression

$$P_{\text{esc}} = \frac{1}{1 + \frac{\Sigma_F \bar{l}_F}{A}}, \quad (16)$$

where A = Levine factor (fuel-rod-geometry dependent). Equation (16) preserves the convenient form of the Wigner rational expression at the two extreme limits and, in addition, it provides good values of P_{esc} for intermediate-size lumps. Incidentally, for A equal to unity, Eq. (16) reduces to Wigner's approximation.

For cylindrical rods, Otter²⁴ has found that the energy-independent value of 1.35 for A works quite well for a wide range of fuel-rod radii. When Eq. (16) is substituted into Eq. (15), the resulting expression for P_{esc}^* is

$$P_{\text{esc}}^* = \frac{1}{1 + \frac{\Sigma_F}{\Sigma_e}}, \quad (17)$$

where the effective cross section Σ_e is given by

$$\Sigma_e = \frac{A(1 - C)}{\bar{l}_F[1 + C(A - 1)]}. \quad (18)$$

The advantage of the rational form of Eq. (17) is the equivalence between the given heterogeneous system and a corresponding homogenized system for which the moderator cross section equals the moderator cross section in the fuel rod

of the heterogeneous system plus the effective cross section Σ_e .^{25,26} This implies that fuel-rod heterogeneity corrections to homogeneous cross sections can be made by adding Σ_e to the fuel-rod moderator cross section and treating the reactor system as homogeneous.

This formalism has been discussed in detail elsewhere^{25,26,27} and has been included in a modified version of the LDX code.

B. Second Method

The second method is in a way an extension of the fuel-rod heterogeneity correction and accounts for both levels of heterogeneity by means of collision probabilities and rational approximations.

From results of the first method, we have found that corrections associated with the "fine" (grain) heterogeneity in HTGR rods of the type under consideration (containing low-volume fractions of 200- to 500- μm -diam grains) is considerably smaller than the "gross" (fuel-rod) heterogeneity correction. Consequently, it would be possible to extend the rational-approximation collision-probability methods of the "gross" heterogeneity correction in order to account for both levels of heterogeneity. The method is briefly as follows.

Let us first define the following quantities:

P_E^* = neutron escape probability from the fuel in the reactor core,

P_e = escape probability from one grain for neutrons uniformly and isotropically produced in that homogeneous grain,

P_E = escape probability from a homogenized fuel rod for neutrons produced uniformly and isotropically in that fuel rod,

f_0 = volume fraction of the grains in one fuel rod,

P_F = probability that a neutron incident on a fuel rod collides in that fuel rod,

P_M = probability that a neutron leaving a fuel rod collides in the moderator outside that rod,

P_0 = probability that a neutron incident on a fuel grain collides in that grain,

P_1 = probability that a neutron leaving a fuel grain collides in the moderator outside it but inside the fuel rod in which the grain is,

P_E' = neutron escape probability from a fuel rod for neutrons produced in the grains of that fuel rod,

P_{ge} = probability that a neutron from the moderator outside any grain will escape from the fuel rod in which that grain is.

From these definitions, it immediately follows that

$$C = 1 - P_M \quad (19)$$

and

$$C_0 = 1 - P_1 \quad , \quad (20)$$

where

C = Dancoff factor of the regular array of fuel rods in the reactor core,
and

C_0 = Dancoff factor of the grains in a fuel rod, i.e., the probability that a neutron leaving a grain will next collide with another grain of the same fuel rod.

From reciprocity theorems,²⁰ it also follows that

$$P_F = \Sigma_F \bar{\ell}_F P_E \quad (21)$$

and

$$P_0 = \Sigma_0 \bar{\ell}_0 P_e \quad , \quad (22)$$

where

Σ_0 = macroscopic fuel-grain cross section,

$\bar{\ell}_0 = \frac{4V_0}{S_0}$ = mean chord length of a grain of volume V_0 and surface area S_0 ;
for a spherical grain of radius R , $\bar{\ell}_0 = (4/3) R$.

The overall neutron escape probability is given by:

$$P_E^* = P_E' [P_M + (1 - P_M)(1 - P_F)P_M + \dots] = P_E' \frac{P_M}{1 - (1 - P_M)(1 - P_F)} \quad , \quad (23)$$

or, combining Eqs. (19), (21), and (23), one obtains

$$P_E^* = P_E' \frac{1 - C}{1 - C(1 - \Sigma_F \bar{\ell}_F P_E)} \quad . \quad (24)$$

The rational approximations for P_E and P_e are

$$P_E = \frac{1}{1 + \frac{\Sigma_F \bar{\ell}_F}{A}} \quad (25)$$

and

$$P_e = \frac{1}{1 + \frac{\Sigma_0 \bar{\ell}_0}{a}} \quad , \quad (26)$$

where A is the rod-geometry-dependent Levine factor¹⁵ with the recommended value²⁴ of 1.35 for cylindrical rods. Parameter "a" can be obtained by "rationalizing" Eq. (10) to give

$$P_e^{\text{sph}} \approx \frac{1}{1 + \frac{9}{16} \Sigma_0 \bar{\ell}_0} \quad , \quad (27)$$

i.e., assigning the value of 16/9 to the Levine-like parameter "a."

We can evaluate P_E' from the series: .

$$P_E' = P_e [P_1 P_{ge} + (1-P_1)(1-P_0)P_1 P_{ge} + \dots] = P_e \frac{P_1 P_{ge}}{1 - (1-P_1)(1-P_0)} \quad , \quad (28)$$

which, after combining Eqs. (20), (26), (22), and (28), yields

$$P_E' = \frac{P_{ge}}{1 + \Sigma_0 \bar{\ell}_0 \left[\frac{1}{a} + \frac{C_0}{1-C_0} \right]} \quad . \quad (29)$$

If we now treat the grains-in-the-fuel-rod configuration as a perturbation of the homogeneous rod model, we can replace Eq. (29) by the approximate expression

$$P_E' \approx \frac{P_E}{1 + \Sigma_F \bar{\ell}_0 \left[\frac{1}{a} + \frac{C_0}{1-C_0} \right]} \quad . \quad (30)$$

Equations (24), (25), and (30) can be combined to give:

$$P_E^* = \frac{1}{\left[1 + \Sigma_F \bar{\ell}_0 \left(\frac{1}{a} + \frac{C_0}{1 - C_0}\right)\right] \left[1 + \Sigma_F \bar{\ell}_F \left(\frac{1}{A} + \frac{C}{1 - C}\right)\right]}, \quad (31)$$

which after neglecting second-order terms yields

$$P_E^* = \frac{1}{1 + \Sigma_F \bar{\ell}_F \left(\frac{1}{A^*} + \frac{C}{1 - C}\right)}, \quad (32)$$

where

$$\frac{1}{A^*} = \frac{1}{A} + \frac{\bar{\ell}_0}{\bar{\ell}_F} \left(\frac{1}{a} + \frac{C_0}{1 - C_0}\right). \quad (33)$$

Equation (32) preserves the rational form of Eq. (16) and corrects for both levels of heterogeneity provided that the Levine parameter A is replaced by the new grain-dependent parameter A* given by Eq. (33). Equation (33) can be written as

$$P_E^* = \frac{1}{1 + \frac{\sigma_F}{\sigma_{\text{eff}}}}, \quad (34)$$

where

$$\sigma_{\text{eff}} = \frac{1}{N_F \bar{\ell}_F \left(\frac{1}{A} + \frac{C}{1 - C}\right)}. \quad (35)$$

N_F = absorber atomic density in the fuel rod. All the σ 's are microscopic cross sections per absorber atom. The new quantity σ_{eff} can then replace Σ_e/N_F of Eq. (18) in the single-heterogeneity correction discussed in Sec. IV.A.2 to yield double-heterogeneity corrections.

This method can be easily incorporated in codes like MC²-I or 1DX without need of pointwise cross sections as required by the first double-heterogeneity shielding method discussed in Sec. IV.A.1.

A similar space shielding method was developed earlier and is discussed in Ref. 28. The grain Dancoff factor calculation necessary for Eq. (33) is derived in Refs. 28 and 29 and is given by:

$$C_0 = \frac{\Sigma_g}{\Sigma_f} \left\{ 1 - \left[1 + \frac{\Sigma_f \bar{\sigma}_F}{m+1} \right]^{-(m+1)} \right\}, \quad (36)$$

where

$$\Sigma_g = n \bar{\sigma}_0, \quad (37)$$

$$\Sigma_f = \Sigma_g + \Sigma_{\text{mod}}, \quad (38)$$

$$\Sigma_{\text{mod}} = N_1 \sigma_1, \quad (39)$$

and

N_1 = atomic density of fuel-rod moderator outside the grains,

σ_1 = fuel-rod moderator microscopic cross section,

$n = f_0/V_0$ = number of grains per unit volume of the fuel rod,

$\bar{\sigma}_0 = \frac{S_0}{4}$ = average "geometric" cross section of the grains,

$m = 3.58$.

If scattering effects in the fuel grains are considered, parameter "a" should be replaced^{28,29} by group parameter a^* :

$$a^* = \frac{a}{1-q}, \quad (40)$$

where q is the ratio of the self-scattering cross section to the total cross section in a particular group. Scattering effects in fuel grains are generally of relatively small importance for the HTGR rods under consideration.

C. Comparisons and Discussion

The above double-heterogeneity space-shielding methods were used for generating above-thermal few-group ^{232}Th , ^{235}U , and ^{233}U cross sections for a 3000-MW(th) HTGR system with fuel rods containing 500- and 200- μm -diameter ThO_2 and UC_2 grains, respectively, in a graphite matrix. The most affected in the above-thermal region is the ^{232}Th absorption cross section of group 3 (in the group structure of Table I), which incorporates all resolved resonances of Thorium. Table III shows a comparison of the group-3 absorption cross sections at 3 temperatures (300, 800, and 1200 K) as calculated by the first method (Sec. IV. A), by the second method (Sec. IV. B), and by the GA code MICROX (GA results supplied to LASL on magnetic tape). A non-grain-shielded absorption cross section (NGSX) is also included for comparison. The grain-shielding effect is seen to be of the order of 4-5% by comparison with the fuel-rod shielding effect, which was seen to be ~25%. In the thermal region, the space shielding of the ^{233}U and ^{235}U absorption cross sections (^{232}Th is not important in the thermal region) was seen to be considerably less important.

TABLE III
RESOLVED-RESONANCE-GROUP ABSORPTION
CROSS SECTION IN ^{232}Th (b)

<u>Temperature</u> (K)	<u>1st</u> <u>Method</u>	<u>2nd</u> <u>Method</u>	<u>MICROX</u>	<u>NGSX</u>
300	6.58	6.72	6.76	6.95
800	7.82	8.03	8.12	8.28
1200	8.42	8.65	8.78	8.90

APPENDIX A

PETOPES PROGRAM

The purpose of the PETOPES program is to change a PENDF tape to a PENDF shielded tape; that is, to produce a pointwise tape in the ENDF/B format containing grain-shielded cross sections from a pointwise ENDF/B tape originally produced by the MINX¹¹ code. The shielded data can then be used as input to the MINX code to obtain multigroup grain-shielded cross sections.

The grain-shielding technique used in PETOPES is that suggested by Wälti.¹³ Although the theory is discussed in detail in the text, the formulas used in the Wälti treatment are repeated here in a notation mnemonically compatible with that used in the code. Grain shielding may be accounted for by noting that the effective resonant material (e.g., thorium in the HTGR) cross section is given by

$$\sigma_{\text{Th}}^{\text{eff}} = \sigma_{\text{Th}} \frac{V_f}{V_c} \frac{\Gamma(E)}{1 + \frac{V_p}{V_c} \Gamma(E)}, \quad (\text{A-1})$$

where σ_{Th} is the unshielded cross section, V_f , V_p , and V_c are the relative volumes of fuel, particle, and moderator regions, respectively, and $\Gamma(E)$ is the energy-dependent disadvantage factor for the particle relative to the remainder of the fuel element. $\Gamma(E)$ depends on the energy-dependent total and scattering cross sections of the resonant material and on other parameters which are insensitive to energy. $\Gamma(E)$ is given by Wälti as

$$\Gamma(E) = \frac{1 + \frac{V_p}{V_c} Q(1 + \tau_{x,c} W)}{1 + \frac{V_p}{V_c} Q + \tau_{x,p} W}, \quad (\text{A-2})$$

where p refers to the particle region, c refers to the moderator region, ξ 's are

$$\tau_{x,j} = \tau_{t,j} \left[1 - \left(1 - \frac{\xi_j}{\xi_c^{\text{pot}}} \right) \frac{\Sigma_{s,j}}{\Sigma_{t,j}} \right] \quad j = p, c, \quad (\text{A-3})$$

the logarithmic slowing-down decrements for each region, and Σ_s and Σ_t are macroscopic scattering and total cross sections, respectively, for the resonant material in each region. Note that for region c the potential scattering cross section is used to evaluate ξ , so that this quantity is energy independent in the moderator region.

$$\tau_{t,j} = \frac{4V_j}{S_j} \Sigma_{t,j} \quad , \quad j = p, c \quad , \quad (A-4)$$

where S refers to the surface areas of the regions.

$$W = 1 + \tilde{H}_0(\tau_{t,p}) + \tilde{H}_1(\tau_{t,c}) \quad . \quad (A-5)$$

$$\tilde{H}_0(\tau_{t,p}) = \frac{1 - \tilde{P}_0(\tau_{t,p})}{\tau_{t,p} \tilde{P}_0(\tau_{t,p})} \quad . \quad (A-6)$$

$$\tilde{P}_0(\tau_{t,p}) = \frac{3}{8X^3} [2X^2 - 1 + (1 + 2X)e^{-2X}] \quad , \quad X = \frac{3}{4} \tau_{t,p} \quad . \quad (A-7)$$

$$\begin{aligned} \tilde{H}_1(\tau_{t,c}) = \left(\frac{\gamma}{r}\right)^2 & \left\{ (1-r^2)^2 \left(1 + \frac{1}{4} \ln \frac{1+r}{1-r}\right) - \frac{r}{2} (1-r)^2 \right. \\ & \left. + \left(\frac{2}{3r}\right)^2 \left[(1-r^2)^3 - 3(1-r^3)^2 + 2(1-r^3)(1-r^2)^{3/2} \right] \right\} . \end{aligned} \quad (A-8)$$

$$r = R_0/R_1 \quad , \quad (A-9)$$

where R_0 and R_1 are outer radii of regions p and c, respectively.

$$\gamma = \frac{3r^2}{4(1-r^3)} \quad . \quad (A-10)$$

$$Q = \frac{\xi_p^{\text{pot}} \Sigma_p^{\text{pot}}}{\xi_c^{\text{pot}} \Sigma_c^{\text{pot}}} \quad . \quad (A-11)$$

Also, the cross-section weighted logarithmic decrements for the mixtures in each region are given by

$$\xi_j = \frac{\sum_k \xi_k N_k^j \sigma_{sk}}{\sum_k N_k^j \sigma_{sk}} \quad , \quad (A-12)$$

where the N_k are the concentration and σ_{sk} the scattering cross sections for isotopic constituents of the regions.

The basic input to the PETOPES code is a PENDF file output by the MINX code. This file usually consists of the cross-section data for a particular nuclide (e.g., ^{232}Th) given for several temperatures. The object of the PETOPES code is to calculate a grain-shielding factor (Eq. A-1) at each energy point in the PENDF file, multiply this factor by the cross section at the given energy, and prepare a new file of the grain-shielded cross sections. This is done for every temperature on the tape. If there is more than one nuclide in a mixture contributing to the grain shielding, a preparatory routine, DBLSHLD, is called which prepares a cross-section file used in calculating the shielding factors according to the formula:

$$\sigma_{\text{eff}} = \sum_{i=1}^n N_i \sigma_i \quad , \quad (A-13)$$

where σ_{eff} is the effective cross section for calculating the self-shielding factor at a particular energy point; n the number of nuclides in the mixture contributing to the self-shielding; N_i the fraction of the i -th nuclide in the mixture, and σ_i the cross section of the i -th nuclide at the energy point in question.

In the data input to the PETOPES code, only the cross-section data for the material for which grain-shielded cross sections are being prepared are assumed to be energy dependent. Total and potential cross sections as well as logarithmic decrements for other materials in the mixtures are assumed to be energy independent. Other input parameters are the radii of the particle and moderator regions and the concentrations of the constituents of particle and surrounding moderator regions. Also the energy range over which the grain shielding is applied is specified. Input specifications are given in Table A-I.

TABLE A-I

PETOPES INPUT SPECIFICATIONS

<u>Card No.</u>	<u>Format</u>	<u>Variable</u>	<u>Comment</u>
1	6A10	A(I)	Title card.
2	6E11.4	RADP	Radius of particle region.
		RADC	Radius of moderator region.
		EMAX	Upper energy bound of resonance region.
		EMIN	Lower energy bound of resonance region.
3	6I11	NMP	No. of materials in particle region.
		NMC	No. of materials in moderator region.
		NOQCAL	Obsolete.
4	6E11.4	PSIP(I)	NMP values of ξ_i for the materials in particle region. Note I=1 is always material for which grain-shielded cross sections are being produced, e.g., Th.
5	6E11.4	PSIC(I)	NMC values of ξ_i for the materials in moderator region. Note I=1 is always for the moderating material, e.g., c.
6	6E11.4	CONP(I)	NMP concentrations for the materials in the particle region. Order same as for PSIP.
7	6E11.4	CONC(I)	NMC concentrations for the materials in the moderator region. Order same as for PSIC.
8	6E11.4	XSP(I),XP(I)	NMP values for total and potential cross sections for materials in particle region. Order same as for PSIP but XSP(I) and XP(I), for the grain-shielded material, are not used because the energy-dependent cross sections are read from input tape.
9	6E11.4	XSC(I),XC(I)	NMC values for total and potential cross sections for materials in the moderator region. Order same as for PSIC.

Comparison of $\Gamma(E)$ as computed by the PETOPES with a calculation of Wälti's¹³ for the 21.8 and 23.5 eV ^{232}Th resonances for ThC_2 particles is shown in Fig. A-1. The agreement is good and differences are attributed to the fact that a different evaluation for ^{232}Th (ENDF/B-III) was used in the PETOPES code from that used by Wälti. This is evident from the fact that the resonances occur at slightly different energies. Figure A-2 shows the variation of $\Gamma(E)$ with temperature for the same two resonances.

A listing of the PETOPES code is given at the end of this appendix. In addition to the grain-shielded file output by the code, printed output includes the input and a limited number of grain-shielding factors and values of $\Gamma(E)$ for each temperature. Plots are also made of these for the various temperatures.

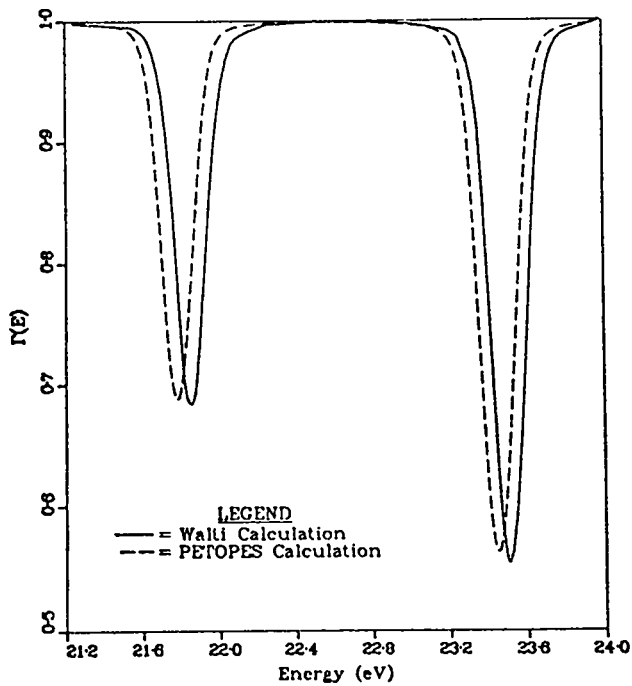


Fig. A-1.
Comparison of Wälti and PETOPES calculations for $\Gamma(E)$ for the 21.8- and 23.5-eV resonances of ^{232}Th at 300 K.

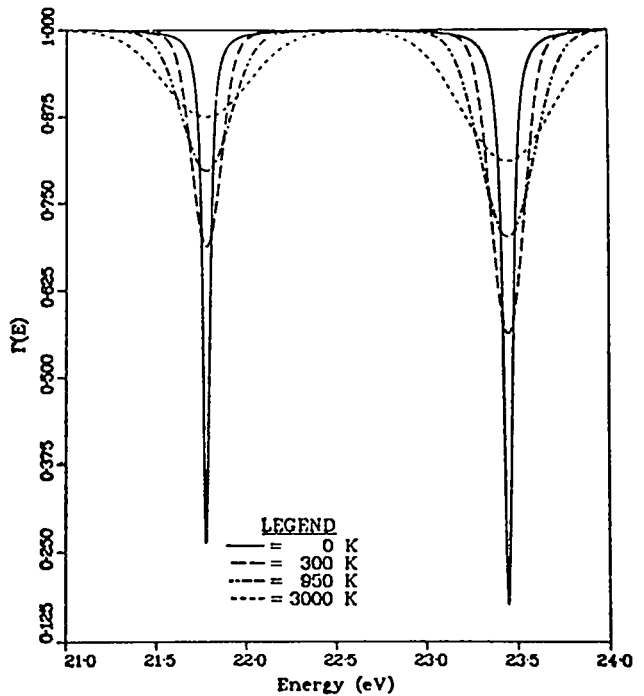


Fig. A-2.
 $\Gamma(E)$ for 0, 300, 950, and 3000 K for the 21.8- and 23.5-eV resonances of ^{232}Th .

```

PROGRAM PETOPES (INP,OUT,FSET5=INP,FSET6=OUT,FSET10,FSET11,FSET12,PETOP 1
1 FILM,FSET9) PETOP 2
C PURPOSE OF PROGRAM - TO CONVERT A PENDF TAPE TO SHIELDED PENDF. PETOP 3
C PENDF TO PENDF-SHIELDED. PETOP 4
C -- -- -- PETOP 5
LCM/XSECTT/XT(60000),YT(60000),NPTT PETOP 6
LCM/XSECTE/XE(60000),YE(60000),NPEE PETOP 7
COMMON/CONS/RADC,RADP,VOLC,VOLP,SURC,SURP,PSIP(10),PSIC(10),EMAX, PETOP 8
1 EMIN,MT PETOP 9
COMMON/CALC/H1TAU,VOLF,Q,SEEP,SEEC,SIGPSP,SIGPSC,TAUTC,TAUXC PETOP 10
COMMON/CON1/CONP(10),CONC(10),XSC(9),XSP(9),NMP,NMC,XP(9),XC(9) PETOP 11
COMMON/PIOTS/ENG(5000),FAX(5000),GAMX(5000),NX,TITL(5),XLR(5), PETOP 12
1 YLB(5) PETOP 13
DIMENSION F(10),S(10),J(10),A(8),HOL(7),X(10),Y(10) PETOP 14
4 READ (11,15) (A(I),I=1,6),ANEXT1,MCHECK,ANEXT2 PETOP 15
5 WRITE (10,15) (A(I),I=1,6),ANEXT1,MCHECK,ANEXT2 PETOP 16
IF (MCHECK.EQ.4H -1) GO TO 6 PETOP 17
GO TO 4 PETOP 18
6 END FILE 10 PETOP 19
REWIND 10 PETOP 20
REWIND 11 PETOP 21
15 FORMAT (6A10,A6,A4,A10) PETOP 22
C INPUT DEFINITIONS- PETOP 23
C RADP - RADIUS OF PARTICLE,E.G. THORIUM CORE OF THORIUM COATED PETOP 24
C PARTICLE IN HTGR PETOP 25
C RADC - RADIUS OF EFFECTIVE SPHERICAL SHELL,E.G. RADIUS OF EFFECTIVE PETOP 26
C MEDIA SURROUNDING THORIUM CORE IN HTGR FUEL ELEMENT. PETOP 27
C VOLP-PARTICLE VOLUME CORRESPONDING TO RADP. PETOP 28
C VOLC-VOLUME CORRESPONDING TO MEDIA SURROUNDING PARTICLE REGION. PETOP 29
C SURP-SURFACE AREA OF PARTICLE. PETOP 30
C SURC-SURFACE AREA OF SURROUNDING MEDIA. PETOP 31
C PSIP- LOG-DEC (MT252) FOR MATERIALS IN PARTICLE REGION,E.G. PETOP 32
C FOR THORIUM PSIP= 0.008669. PETOP 33
C PSIC- LOG-DEC (MT252) FOR MATERIALS OUTSIDE,PARTICLE REGION. FOR PETOP 34
C CARRON,PSIC= 0.1589. PETOP 35
C NMP=NO OF MATS IN PARTICLE REGION. PETOP 36
C NMC=NO OF MATS IN OUTER REGION. PETOP 37
C CONP-ATOMS/CC OF MATS IN PARTICLE REGION. CONP(1) IS FOR THORIUM PETOP 38
C CONC-ATOMS/CC OF MATS OUTSIDE PARTICLE REGION. PETOP 39
C XSP,XP-TOT,POT XSEC FOR MATERIALS WITH CONSTANT XSEC IN PARTICLE PETOP 40
C REGION. XSP(1),XP(1),ARE FOR THORIUM-COMPUTED IN GRANSHL. PETOP 41
C XSC,XC-TOT,POT XSEC FOR MATERIALS OUTSIDE PARTICLE REGION. PETOP 42
C EMAX-ENERGY BOUNDING RESONANCE REGION FOR PARTICLE SHIELDING,E.G. PETOP 43
C EMAX=4.0 KEV FOR THORIUM. PETOP 44
C EMIN- LOWER BOUND OF RESONANCE REGION,E.G.,EMIN=21EV FOR TH-232. PETOP 45
C NOQCAL = 0 FOR FERTILE MATS,F.G. THORIUM IN RES. REGION. PETOP 46
C = 1 FOR FISSILE MATS,F.G. U-235 AND U-233 IN THERMAL PETOP 47
C REGION ONLY.(NOTE THERMAL REGION MUST BE RUN SOLO PETOP 48
C BECAUSE OF THIS) PETOP 49
READ (5,15) (A(I),I=1,6) PETOP 50
WRITE (6,15) (A(I),I=1,6) PETOP 51
READ (5,18) RADP,RADC,EMAX,EMIN PETOP 52
7 FORMAT (1H0,* INPUT*//* RADP = *1PE12.5,* RADC = *1PE12.5, PETOP 53
1 * EMAX = *1PE12.5,* EMIN = *1PE12.5,* NOQCAL = *I3) PETOP 54
18 FORMAT (6F11,4) PETOP 55
READ (5,19) NMP,NMC,NOQCAL PETOP 56
WRITE (6,7) RADP,RADC,EMAX,EMIN,NOQCAL PETOP 57
19 FORMAT (6I11) PETOP 58
READ (5,18) (PSIP(I),I=1,NMP) PETOP 59
READ (5,18) (PSIC(I),I=1,NMC) PETOP 60
READ (5,18) (CONP(I),I=1,NMP) PETOP 61

```

	READ (5,18) (CONC(I),I=1,NMC)	PETOP 63
	READ (5,18) (XSP(I),XP(I),I=1,NMP)	PETOP 64
	READ (5,18) (XSC(I),XC(I),I=1,NMC)	PETOP 65
C		PETOP 66
C	THESE CONSTANTS ARE NEEDED IN SUBROUTINE GRANSHL.	PETOP 67
		PETOP 68
	VOLP=4./3.*3.14159*RADP**3	PETOP 69
	VOLC=4./3.*3.14159*RADC**3 -VOLP	PETOP 70
	SURP=4.*3.14159*RADP**2	PETOP 71
	SURC=SURP	PETOP 72
	R=RADP/RADC	PETOP 73
	GAM=3*R**2/(4.0*(1.0-R**3))	PETOP 74
	TRM1=(1.0-R**2)**2*(1.0+0.25*ALOG((1.0+R)/(1.0-R)))	PETOP 75
1	0.5*R*(1.0-R)**2	PETOP 76
	TRM2=(2.0/(3.0*R))**2	PETOP 77
	TRM3=(1.0-R**2)**3-3.0*(1.0-R**3)**2+2.0*(1.0-R**3)*(1.0-R**2)**1	PETOP 78
1 5		PETOP 79
	HITAU=(GAM/R)**2*(TRM1+TRM2+TRM3)	PETOP 80
	VOLF=VOLC+VOLP	PETOP 81
C		PETOP 82
C	CALCULATE Q.	PETOP 83
C		PETOP 84
	SEENUM=0.	PETOP 85
	SIGPSP=0.	PETOP 86
	DO 130 I=1,NMP	PETOP 87
	SEENUM=PSIP(I)*CONP(I)*XP(I)+SEENUM	PETOP 88
130	SIGPSP=CONP(I)*XP(I)+SIGPSP	PETOP 89
	CONTINUE	PETOP 90
	Q=SEENUM	PETOP 91
	SEEP=SEENUM/SIGPSP	PETOP 92
	SEENUM=0.	PETOP 93
	SIGPSC=0.	PETOP 94
	SIGTC=0.	PETOP 95
	DO 140 I=1,NMC	PETOP 96
	SEENUM=PSIC(I)*CONC(I)*XC(I)+SEENUM	PETOP 97
	SIGPSC=CONC(I)*XC(I)+SIGPSC	PETOP 98
	SIGTC=CONC(I)*XSC(I)+SIGTC	PETOP 99
140	CONTINUE	PETOP100
	SEEC=SEENUM/SIGPSC	PETOP101
	Q=Q/SEENUM	PETOP102
	TAUTC=4.0*VOLC*SIGTC/SURC	PETOP103
	TAUXC=TAUTC	PETOP104
C		PETOP105
C	END OF Q CALCULATION	PETOP106
C		PETOP107
	TITL(1)=10HGAMMA PLOT	PETOP108
	TITL(2)=10H TO COMPAR	PETOP109
	TITL(3)=10HE WITH OTH	PETOP110
	TITL(4)=10HR METHODS.	PETOP111
	XLB(1)=10HENERGY IN	PETOP112
	XLB(2)=10HF.V. UNITS	PETOP113
10	READ (11,20) (HOL(I),I=1,7),MAT,MF,MT,NSEQ	PETOP114
	HOL(1)=10H THIS TAP	PETOP115
	HOL(2)=10HF HAS BEEN	PETOP116
	HOL(3)=10H CHANGED T	PETOP117
	HOL(4)=10H A PENDF=	PETOP118
	HOL(5)=10HSHIELDED F	PETOP119
	HOL(6)=10HILE.	PETOP120
20	FORMAT (A10,A6,I4,I2,I3,I5)	PETOP121
	WRITE (12,20) (HOL(I),I=1,7),MAT,MF,MT,NSEQ	PETOP122
	READ (10,20) DUM	PETOP123
	READ (10,80) ZA,AWR	PETOP124
	CALL STORXS	PETOP125

```

NX=0
PRINT 2020,MAT
2020 FORMAT (1H1,*WELL,WE MADE IT OUT OF STORXS ONCE, MAT=*I4)
30 READ (11,20) (HOL(I),I=1,7),MAT,MF,MT,NSEQ
IF (MAT.EQ.0) CALL STORXS
IF (MAT.EQ.0) NX=0
2030 FORMAT (1H,*WE ARE LOOPING NOW, MAT=*I4)
WRITE (12,20) (HOL(I),I=1,7),MAT,MF,MT,NSEQ
IF (MAT.EQ.-1) GO TO 2000
IF (MF.NE.3) GO TO 30
IF (MT.EQ.1) GO TO 31
IF (MT.EQ.2) GO TO 31
IF (MT.EQ.3) GO TO 31
IF (MT.EQ.18) GO TO 31
IF (MT.EQ.102) GO TO 31
GO TO 30
31 CONTINUE
MTXX=MT
READ (11,40) C1,C2,N1,N2,N3,N4,MAT,MF,MT,NSEQ
CALL CXFP (C1,F(1),S(1),J(1))
CALL CXFP (C2,F(2),S(2),J(2))
WRITE (12,50) (F(I),S(I),J(I),I=1,2),N1,N2,N3,N4,MAT,MF,MT,NSEQ
40 FORMAT (1P2E11.4,4I11,I4,I2,I3,I5)
50 FORMAT (2(F8.5,A1,I2),4I11,I4,I2,I3,I5)
READ (11,60) NPT,INT,N0,N0,N0,N0,N0,MAT,MF,MT,NSEQ
WRITE (12,60) NPT,INT,N0,N0,N0,N0,N0,MAT,MF,MT,NSEQ
WRITE (6,100) C1,MAT,MT
60 FORMAT (6I11,I4,I2,I3,I5)
NN1=1
70 NN2=NN1+2
READ (11,80) (X(I),Y(I),I=1,3),MAT,MF,MT,NSEQ
80 FORMAT (1P6E11.4,I4,I2,I3,I5)
LOOP=0
DO 85 I=1,3
E=X(I)
CALL GRANSHL (E,FACT)
Y(I)=Y(I)*FACT
LOOP=LOOP+1
85 CONTINUE
CALL CXFP (X(1),F(1),S(1),J(1))
CALL CXFP (Y(1),F(2),S(2),J(2))
CALL CXFP (X(2),F(3),S(3),J(3))
CALL CXFP (Y(2),F(4),S(4),J(4))
CALL CXFP (X(3),F(5),S(5),J(5))
CALL CXFP (Y(3),F(6),S(6),J(6))
WRITE (12,90) (F(I),S(I),J(I),I=1,6),MAT,MF,MT,NSEQ
90 FORMAT (6(F8.5,A1,I2),I4,I2,I3,I5)
95 FORMAT (* M = *I6,* E = *1PE12.5,* FACT = *1PE12.5)
100 FORMAT (1H1,* TEMPERATURE = *1PE12.5,* MAT = *I4,* MT = *I3)
NN1=NN2+1
IF (NN1.LE.N4) GO TO 70
READ (11,20) (HOL(I),I=1,7),MAT,MF,MT,NSEQ
WRITE (12,20) (HOL(I),I=1,7),MAT,MF,MT,NSEQ
IF (MTXX.GT.1) GO TO 30
WRITE (6,200) NX,C1
200 FORMAT (1H1,* NX = *I6,* FOR TEMP = *1PE12.5//7X,*ENERGY*,15X,
1 *FACT*13X,*GAMMA*)
WRITE (6,210) (ENG(N),FAX(N),GAMX(N),N=1,NX)
210 FORMAT (1P3E18.5)
WRITE (9) NX,(ENG(N),FAX(N),GAMX(N),N=1,NX),NX
GO TO 30
2000 WRITE (6,2010) MAT
2010 FORMAT (1H1,* PROCESSING COMPLETE. MAT = *I4)

```

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PETOP126
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PETOP185
PETOP186
PETOP187
PETOP188

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CALL GFPLT
END

PETOP189
PETOP190

```

SUBROUTINE STORXS
STORE TOTAL AND ELASTIC XSEC FOR THORIUM.
LCM/XSECTT/XT(60000),YT(60000),NPTT
LCM/XSECTE/XE(60000),YE(60000),NPEE
DIMENSION A(7)
10 READ (10,20) (A(I),I=1,7),MAT,MF,MT,NSEQ
20 FORMAT (6A10,A6,I4,I2,I3,I5)
   IF (MAT.EQ.-1) GO TO 2000
   IF (MF.GT.3) GO TO 300
   IF (MF.LT.3) GO TO 10
   IF (MT.EQ.1) GO TO 100
   IF (MT.EQ.2) GO TO 200
   GO TO 10
100 READ (10,30) NPTT
    PRINT 2020,NPTT
    READ (10,20) (A(I),I=1,7)
    IF (NPTT.GT.60000) GO TO 130
    NPTTS=NPTT
105 READ (10,40) (XT(I),YT(I),I=1,NPTTS)
    PRINT 40,(XT(I),YT(I),I=1,99)
    PRINT 2040,NPTTS,XT(NPTTS),YT(NPTTS)
110 IF (NPTT.EQ.NPTTS) GO TO 10
120 IF (MT.EQ.0) GO TO 10
    READ (10,20) (A(I),I=1,7),MAT,MF,MT
    GO TO 120
130 NPTTS=60000
    GO TO 105
200 READ (10,30) NPEE
    PRINT 2030,NPEE
    READ (10,20) (A(I),I=1,7)
    IF (NPEE.GT.60000) GO TO 230
    NPEES=NPEE
205 READ (10,40) (XE(I),YE(I),I=1,NPEES)
    PRINT 40,(XE(I),YE(I),I=1,99)
    PRINT 2050,NPEES,XE(NPEES),YE(NPEES)
210 IF (NPEE.EQ.NPEES) GO TO 10
220 IF (MT.EQ.0) GO TO 10
    READ (10,20) (A(I),I=1,7),MAT,MF,MT
    GO TO 220
230 NPEES=60000
    GO TO 205
300 READ (10,20) (A(I),I=1,7),MAT,MF,MT,NSEQ
    IF (MAT.EQ.-1) GO TO 2000
    IF (MAT.NE.0) GO TO 300
    RETURN
30 FORMAT (55X,I11)
40 FORMAT (6E11,4)
2000 WRITE(6,2010) MAT
2010 FORMAT (1H1,* SORRY TAPE IS OUT OF TEMPS. MAT=*I4)
2020 FORMAT (1H1,10X,*XT,YT TARLF*10X,*NPTT=*I11)
2030 FORMAT (1H0,10X,*XE,YE TARLF*10X,*NPEE=*I11)
2040 FORMAT (1H ,*NPTTS=*I6,4X,*XT(NPTTS)=*E11,4,4X,*YT(NPTTS)=*E11,4)
2050 FORMAT (1H ,*NPEES=*I6,4X,*XE(NPEES)=*E11,4,4X,*YE(NPEES)=*E11,4)
RETURN
END
```

STORX 1
STORX 2
STORX 3
STORX 4
STORX 5
STORX 6
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STORX 8
STORX 9
STORX 10
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STORX 54
STORX 55

	SUBROUTINE GRANSHL (E,FACT)	GRANS 1
C	PURPOSE = TO CALCULATE SHIELDING FACTOR FOR TWO REGION PARTICLE.	GRANS 2
C		GRANS 3
	LCM/XSECT/XT(60000),YT(60000),NPTT	GRANS 4
	LCM/XSECTE/XE(60000),YE(60000),NPEE	GRANS 5
	COMMON/CONS/RAUC,RADP,VOLC,VOLP,SURC,SURP,PSIP(10),PSIC(10),EMAX,	GRANS 6
1	EMIN,MT	GRANS 7
	COMMON/CALC/H1TAU,VOLF,Q,SEEP,SEEC,SIGPSP,SIGPSC,TAUTC,TAUXC	GRANS 8
	COMMON/CON1/CONP(10),CONC(10),XSC(9),XSP(9),NMP,NMC,XP(9),XC(9)	GRANS 9
	COMMON/PLOTS/ENG(5000),FAX(5000),GAMX(5000),NX,TITL(5),XLB(5),	GRANS 10
1	YLB(5)	GRANS 11
C		GRANS 12
C	CONDITIONAL RETURNS	GRANS 13
C		GRANS 14
	IF (E.GT.1.0E-10) GO TO 10	GRANS 15
	E=0.	GRANS 16
	FACT=0.	GRANS 17
	RETURN	GRANS 18
10	CONTINUE	GRANS 19
	IF (E.LT.EMAX) GO TO 20	GRANS 20
	FACT=1.0	GRANS 21
	RETURN	GRANS 22
20	CONTINUE	GRANS 23
	QQ=Q	GRANS 24
	IF (E.LT.EMIN) QQ=0.	GRANS 25
C		GRANS 26
C	FIND TOTAL AND ELASTIC CROSS SECTIONS CORRESPONDING TO ENERGY E.	GRANS 27
C		GRANS 28
	CALL LOCT1(E,ILK,ILOT)	GRANS 29
	IHIT=ILOT+1	GRANS 30
	CALL LOCT2(E,ILK,ILOE)	GRANS 31
	IHIE=ILOE+1	GRANS 32
	DXTL=XT(ILOT)	GRANS 33
	DYTL=YT(ILOT)	GRANS 34
	DXTH=XT(IHIT)	GRANS 35
	DYTH=YT(IHIT)	GRANS 36
	CALL TERP1 (DXTL,DYTL,DXTH,DYTH,E,ST,2,1)	GRANS 37
	DXEL=XE(ILOE)	GRANS 38
	DYEL=YE(ILOE)	GRANS 39
	DXEH=XE(IHIE)	GRANS 40
	DYEH=YE(IHIE)	GRANS 41
	CALL TERP1 (DXEL,DYEL,DXEH,DYEH,E,SE,2,2)	GRANS 42
C		GRANS 43
	XP(1)=SE	GRANS 44
	XSP(1)=ST	GRANS 45
	SIGPSP1=0.	GRANS 46
	SIGTP=0.	GRANS 47
	SEENUM=0.	GRANS 48
	DO 30 I=1,NMP	GRANS 49
	SEENUM=PSIP(I)*CONP(I)*XP(I)+SEENUM	GRANS 50
	SIGPSP1=CONP(I)*XP(I)+SIGPSP1	GRANS 51
	SIGTP=CONP(I)*XSP(I)+SIGTP	GRANS 52
30	CONTINUE	GRANS 53
	SEEP1=SEENUM/SIGPSP1	GRANS 54
	TAUTP=4.0*VOLP*SIGTP/SURP	GRANS 55
	TAUXP=TAUTP*(1.0-(1.0-SEEP1/SEEC)*SIGPSP1/SIGTP)	GRANS 56
	X=3.0*TAUTP/4.0	GRANS 57
	P0TAU=3.0/(8.0*X**3)*(2.0*X**2-1.0*(1.0+2.0*X)*EXP(-2.0*X))	GRANS 58
	H0TAU=(1.0-P0TAU)/(TAUTP*P0TAU)-1.0	GRANS 59
	W=1.0*H0TAU*H1TAU	GRANS 60
	RHOQ=VOLP/VOLC*QQ	GRANS 61
	UPPER=1.0*RHOQ*(1.0+TAUXC*W)	GRANS 62

UNDER=1.0+RHOQ*TAUXP*W	GRANS 63
GAMMA=UPPER/UNDER	GRANS 64
FACT=VOLP/VOLC*(GAMMA/(1.0+VOLP/VOLC*GAMMA))	GRANS 65
IF (FACT.GT.0.999) GO TO 40	GRANS 66
IF (MT.GT.1) GO TO 40	GRANS 67
IF (E.LT.EMIN) GO TO 40	GRANS 68
IF(NX.GT.1200) GO TO 40	GRANS 69
NX=NX+1	GRANS 70
ENG(NX)=E	GRANS 71
FAX(NX)=FACT	GRANS 72
GAMX(NX)=GAMMA	GRANS 73
40 CONTINUE	GRANS 74
45 FORMAT (IH0,* PLOTS GO ONLY TO *1PE12.5,* E.V.*)	GRANS 75
RETURN	GRANS 76
END	GRANS 77

SUBROUTINE CXFP(X,F,S,N)		CXFP 1
C*****		CXFP 2
C* CONVERT X FOR PUNCHING		*CXFP 3
C* X - FLOATING POINT NUMBER = F*10.0**N		*CXFP 4
C* F - 0.999995 LE F !.T 9.999995		*CXFP 5
C* S - SIGN (HOLLERITH + OR -) OF EXPONENT		CXFP 6
C* N - EXPONENT		*CXFP 7
C*****		CXFP 8
DATA SP/IH*/ ,SM/IH*/		CXFP 9
IF(X.NE.0.0) GO TO 10		CXFP 10
F=0.0		CXFP 11
S=SP		CXFP 12
N=0		CXFP 13
RETURN		CXFP 14
10 N=ALOG10(ABS(X))		CXFP 15
IF (ABS(X)-1.0) 40,20,20		CXFP 16
20 F=X/10.0**N		CXFP 17
S=SP		CXFP 18
IF(ABS(F)-9.999995) 70,30,30		CXFP 19
30 F=F/10.0		CXFP 20
N=N+1		CXFP 21
GO TO 70		CXFP 22
40 N=1-N		CXFP 23
F=X*10.0**N		CXFP 24
S=SM		CXFP 25
IF(ABS(F)-9.999995) 70,50,50		CXFP 26
50 F=F/10.0		CXFP 27
N=N-1		CXFP 28
IF(N) 60,60,70		CXFP 29
60 S=SP		CXFP 30
70 CONTINUE		CXFP 31
RETURN		CXFP 32
END		CXFP 33

	SUBROUTINE TERP1 (X1,Y1,X2,Y2,X,Y,I,NERR)	TERP1 1
C	====INTERPOLATE ONE PT.====	TERP1 2
C	(X1,Y1) AND (X2,Y2) ARE END PTS. OF THE LINE	TERP1 3
C	(X,Y) IS INTERPOLATED POINT	TERP1 4
C	I=INTERPOLATION CODE	TERP1 5
C	NOTE - IF A NEGATIVE OR ZERO ARGUMENT OF A LOG IS DETECTED, THE	TERP1 6
C	INTERPOLATION IS AUTOMATICALLY CHANGED FROM LOG TO LINEAR.	TERP1 7
C	ERROR STOPS - 301 (X1=X2,DISCONTINUITY)	TERP1 8
C	302 (INTERPOLATION CODE IS OUT OF RANGE)	TERP1 9
C	303 (ZERO OR NEGATIVE ARGUMENT FOR INTERPOLATED PT.)	TERP1 10
	5 XA=X1	TERP1 11
	YA=Y1	TERP1 12
	XB=X2	TERP1 13
	YB=Y2	TERP1 14
	XP=X	TERP1 15
	II=I	TERP1 16
	IF ((XB-XA).GT.1.E-10) GO TO 7	TERP1 17
	IF (X.EQ.XA) Y=YA	TERP1 18
	PRINT 6,XA,YA,XB,YB,X,Y,I,NERR	TERP1 19
	6 FORMAT (1H0,* ERROR STOP 301 *1P6E12,5,2I3)	TERP1 20
	RETURN	TERP1 21
	7 CONTINUE	TERP1 22
	IF (II) 10,10,15	TERP1 23
	10 CALL ERROR (302)	TERP1 24
	15 IF (II-5) 20,20,10	TERP1 25
	20 GO TO (25,30,35,60,75), II	TERP1 26
	25 YP=YA	TERP1 27
	IF (XP.EQ.XB) YP=YB	TERP1 28
	GO TO 105	TERP1 29
	30 YP=YA+(XP-XA)*(YB-YA)/(XB-XA)	TERP1 30
	GO TO 105	TERP1 31
	35 IF (XA) 30,30,40	TERP1 32
	40 IF (XB) 30,30,45	TERP1 33
	45 IF (XP) 50,50,55	TERP1 34
	50 CALL ERROR (303)	TERP1 35
	55 YP=YA+ALOG(XP/XA)*(YB-YA)/ALOG(XB/XA)	TERP1 36
	GO TO 105	TERP1 37
	60 IF (YA) 30,30,65	TERP1 38
	65 IF (YB) 30,30,70	TERP1 39
	70 YP=YA*EXP((XP-XA)*ALOG(YB/YA)/(XB-XA))	TERP1 40
	GO TO 105	TERP1 41
	75 IF (YA) 35,35,80	TERP1 42
	80 IF (YB) 35,35,85	TERP1 43
	85 IF (XA) 70,70,90	TERP1 44
	90 IF (XB) 70,70,95	TERP1 45
	95 IF (XP) 50,50,100	TERP1 46
	100 YP=YA*EXP(ALOG(XP/XA)*ALOG(YB/YA)/ALOG(XB/XA))	TERP1 47
	105 Y=YP	TERP1 48
	RETURN	TERP1 49
	END	TERP1 50

	SUBROUTINE TERP1 (X1,Y1,X2,Y2,X,Y,I,NERR)	TERP1 1
	=====INTERPOLATE ONE PT.=====	TERP1 2
C	(X1,Y1) AND (X2,Y2) ARE END PTS. OF THE LINE	TERP1 3
C	(X,Y) IS INTERPOLATED POINT	TERP1 4
C	I=INTERPOLATION CODE	TERP1 5
C	NOTE - IF A NEGATIVE OR ZERO ARGUMENT OF A LOG IS DETECTED, THE	TERP1 6
C	INTERPOLATION IS AUTOMATICALLY CHANGED FROM LOG TO LINEAR.	TERP1 7
C	ERROR STOPS - 301 (X1=X2,DISCONTINUITY)	TERP1 8
C	302 (INTERPOLATION CODE IS OUT OF RANGE)	TERP1 9
C	303 (ZERO OR NEGATIVE ARGUMENT FOR INTERPOLATED PT.)	TERP1 10
	5 XA=X1	TERP1 11
	YA=Y1	TERP1 12
	XB=X2	TERP1 13
	YB=Y2	TERP1 14
	XP=X	TERP1 15
	II=I	TERP1 16
	IF ((XB-XA).GT.1.E-10) GO TO 7	TERP1 17
	IF (X.EQ.XA) Y=YA	TERP1 18
	PRINT 6,XA,YA,XB,YB,X,Y,I,NERR	TERP1 19
	6 FORMAT (1H0,* ERROR STOP 301 *1P6E12,5,2I3)	TERP1 20
	RETURN	TERP1 21
	7 CONTINUE	TERP1 22
	IF (II) 10,10,15	TERP1 23
	10 CALL ERROR (302)	TERP1 24
	15 IF (II-5) 20,20,10	TERP1 25
	20 GO TO (25,30,35,60,75), II	TERP1 26
	25 YP=YA	TERP1 27
	IF (XP.EQ.XB) YP=YB	TERP1 28
	GO TO 105	TERP1 29
	30 YP=YA+(XP-XA)*(YB-YA)/(XB-XA)	TERP1 30
	GO TO 105	TERP1 31
	35 IF (XA) 30,30,40	TERP1 32
	40 IF (XB) 30,30,45	TERP1 33
	45 IF (XP) 50,50,55	TERP1 34
	50 CALL ERROR (303)	TERP1 35
	55 YP=YA+ALOG(XP/XA)*(YB-YA)/ALOG(XB/XA)	TERP1 36
	GO TO 105	TERP1 37
	60 IF (YA) 30,30,65	TERP1 38
	65 IF (YB) 30,30,70	TERP1 39
	70 YP=YA*EXP((XP-XA)*ALOG(YB/YA)/(XB-XA))	TERP1 40
	GO TO 105	TERP1 41
	75 IF (YA) 35,35,80	TERP1 42
	80 IF (YB) 35,35,85	TERP1 43
	85 IF (XA) 70,70,90	TERP1 44
	90 IF (XB) 70,70,95	TERP1 45
	95 IF (XP) 50,50,100	TERP1 46
	100 YP=YA*EXP(ALOG(XP/XA)*ALOG(YB/YA)/ALOG(XB/XA))	TERP1 47
	105 Y=YP	TERP1 48
	RETURN	TERP1 49
	END	TERP1 50

	SUBROUTINE ERROR (N)	ERROR 1
	IOS=9	ERROR 2
	5 PRINT 10,N	ERROR 3
	WRITE(99,10)	ERR\JR 4
	10 FORMAT (11H ERROR STOP,I6)	ERROR 5
	END	ERROR 6

```

SUBROUTINE LOCT1(X,ILO,LOCT)
C BINARY SEARCH ROUTINE WRITTEN BY P. SORAN, MODIFIED 10-30-73
C TO GIVE RESULTS IDENTICAL TO EARLER LOCT ROUTINE.
C THAT IS, FIND X SUCH THAT A(LOCT+1).GT.X.GE.A(LOCT), EXCEPT W
C X IS EQUAL TO A(N). IN THAT CASE, LOCT IS SET TO (N-1).
C WHEN X IS NOT BINNABLE, THAT IS WHEN X IS OUTSIDE THE RANGE OF
C A-VALUES OR IF A CONTAINS ONLY A SINGLE POINT, THE VALUE LOCT=
C IS RETURNED.
LCM/XSECTT/A(60000),YT(60000),N
IF(N.EQ.1) GO TO 3001
IF(X.LT.A(1)) GO TO 3001
IF(X.GT.A(N)) GO TO 3001
IF((A(N-1).EQ.A(N)).AND.(X.EQ.A(N))) GO TO 3001
LOCT=1
IF(A(1).EQ.X) RETURN
ILO=1
ISRCH=N
1000 IF(ISRCH.LE.ILO+1)GO TO 3000
I=(ISRCH+ILO)/2
IF(A(I).LT.X) GO TO 2000
ISRCH=I
GO TO 1000
2000 ILO=I
GO TO 1000
C X HAS BEEN BINNED. CONVERT FROM ISRCH TO LOCT HERE.
3000 IF(X.NE.A(ISRCH)) LOCT=ISRCH-1
IF(X.EQ.A(ISRCH)) LOCT=ISRCH
IF(X.EQ.A(N)) LOCT=N-1
RETURN
3001 WRITE (99,10) LOCT
10 FORMAT (1H,16)
RETURN
END

```

```

SUBROUTINE LOCT2(X,ILO,LOCT)
C BINARY SEARCH ROUTINE WRITTEN BY P. SORAN, MODIFIED 10-30-73
C TO GIVE RESULTS IDENTICAL TO EARLER LOCT ROUTINE.
C THAT IS, FIND X SUCH THAT A(LOCT+1).GT.X.GE.A(LOCT), EXCEPT W
C X IS EQUAL TO A(N). IN THAT CASE, LOCT IS SET TO (N-1).
C WHEN X IS NOT BINNABLE, THAT IS WHEN X IS OUTSIDE THE RANGE OF
C A-VALUES OR IF A CONTAINS ONLY A SINGLE POINT, THE VALUE LOCT=
C IS RETURNED.
LCM/XSECTE/A(60000),XE(60000),N
IF(N.EQ.1) GO TO 3001
IF(X.LT.A(1)) GO TO 3001
IF(X.GT.A(N)) GO TO 3001
IF((A(N-1).EQ.A(N)).AND.(X.EQ.A(N))) GO TO 3001
LOCT=1
IF(A(1).EQ.X) RETURN
ILO=1
ISRCH=N
1000 IF(ISRCH.LE.ILO+1)GO TO 3000
I=(ISRCH+ILO)/2
IF(A(I).LT.X) GO TO 2000
ISRCH=I
GO TO 1000
2000 ILO=I
GO TO 1000
C X HAS BEEN BINNED. CONVERT FROM ISRCH TO LOCT HERE.
3000 IF(X.NE.A(ISRCH)) LOCT=ISRCH-1
IF(X.EQ.A(ISRCH)) LOCT=ISRCH
IF(X.EQ.A(N)) LOCT=N-1
RETURN

```

```

3001 WRITE (99,10) LOCT
10  FORMAT (1H ,I6)
    RETURN
    END

```

```

LOCT2 30
LOCT2 31
LOCT2 32
LOCT2 33

```

<pre> C SUBROUTINE GFPLT SUBROUTINE TO PLOT GAMMA AND FACT FOR REPORT. COMMON/CONS/RADC,RADP,VOLC,VOLP,SURC,SURP,PSIP(10),PSIC(10),EMAX, 1 EMIN,MT COMMON/PLOTS/E(4,1250),F(4,1250),G(4,1250),NX,TITL(5),XLBS(5), 1 YLB(5) DIMENSION YLAB(5),YLAB(5),ZLAB(5),NPT(5),X(5000), 1 Y(5000),Z(5000) REWIND 9 NT=4 DO 10 N=1,NT READ (9) NX,(E(N,I),F(N,I),G(N,I),I=1,NX),NPT(N) 10 CONTINUE NN1=1 X(1)=EMIN Y(1)=1. Z(1)=1. DO 30 N=1,NT JP=NPT(N) DO 20 J=1,JP NN1=NN1+1 X(NN1)=E(N,J) Y(NN1)=F(N,J) Z(NN1)=G(N,J) 20 CONTINUE NN1=NN1+1 X(NN1)=EMAX Y(NN1)=1. Z(NN1)=1. NN1=NN1+1 X(NN1)=EMAX Y(NN1)=0. Z(NN1)=0. NN1=NN1+1 X(NN1)=EMIN Y(NN1)=0. Z(NN1)=0. NN1=NN1+1 X(NN1)=EMIN Y(NN1)=1. Z(NN1)=1. 30 CONTINUE TITL(1)=10H FACT FOR TITL(2)=10H0-300-950. TITL(3)=10H3000 DEG K XLAB(1)=10HEENERGY IN XLAB(2)=10HE.V. UNITS YLAB(1)=10H FACT OF E CALL PLOJR(X,Y,NN1,1,0,0,0,0,1,0,1,0,TITL,30,XLAB,20,YLAB,10) TITL(1)=10HGAMMA FOR YLAB(1)=10HGAMMA OF E CALL PLOJB (X,Z,NN1,1,0,0,0,0,1,0,1,0,TITL,30,XLAB,20,YLAB,10) RETURN END </pre>	<pre> GFPLT 1 GFPLT 2 GFPLT 3 GFPLT 4 GFPLT 5 GFPLT 6 GFPLT 7 GFPLT 8 GFPLT 9 GFPLT 10 GFPLT 11 GFPLT 12 GFPLT 13 GFPLT 14 GFPLT 15 GFPLT 16 GFPLT 17 GFPLT 18 GFPLT 19 GFPLT 20 GFPLT 21 GFPLT 22 GFPLT 23 GFPLT 24 GFPLT 25 GFPLT 26 GFPLT 27 GFPLT 28 GFPLT 29 GFPLT 30 GFPLT 31 GFPLT 32 GFPLT 33 GFPLT 34 GFPLT 35 GFPLT 36 GFPLT 37 GFPLT 38 GFPLT 39 GFPLT 40 GFPLT 41 GFPLT 42 GFPLT 43 GFPLT 44 GFPLT 45 GFPLT 46 GFPLT 47 GFPLT 48 GFPLT 49 GFPLT 50 GFPLT 51 GFPLT 52 GFPLT 53 GFPLT 54 </pre>
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APPENDIX B

ETOGLEN

A code to prepare absorber cross-section input for the GLEN code - ENDF/B to GLEN.

In addition to microscopic cross-section data for the moderating materials, which are supplied by the TOR code, the GLEN thermal multigroup-averaging code also requires pointwise data for elastic scattering, fission, and absorption cross sections for the absorbing materials in the reactor model being calculated. These need be only supplied in the thermal energy range, e.g., up to 2.38 eV for the HTGR but, because of storage limitations in present versions of GLEN, the data must be restricted to fewer than 88 energy cross-section pairs for each reaction. Consequently, some care must be taken in choosing a fine energy grid for a particular problem that adequately reproduces the shapes of the cross sections for all materials present and preserves the resonance integral of the principal constituents.

The purpose of the ETOGLEN code is : a) retrieve pointwise cross-section data from a pointwise ENDF/B file (PENDF), created by the MINX code; b) assist the user in choosing a fine energy grid for a problem by allowing flexible grid input, by providing comparison plots of the selected grid vs the ENDF/B points, and by calculating weighted resonance integrals over specified intervals for data on both the selected grid and the original ENDF/B grid; and c) output absorber cross-section data in the format required by the GLEN code.

The selected grid need not be a subset of the original ENDF/B grid, as the code will interpolate on any given mesh. If a representative spectrum is taken for the weighting function for the resonance integral calculations and broad-group boundaries are taken for the calculational intervals, one obtains the error in the broad-group cross sections incurred by grid selection.

The code calculates an energy grid on the basis of a set of incremental values of lethargy (or velocity increments in another version of ETOGLEN) specified for several energy intervals. To these are added additional points, input by the user, such as peaks and valleys of important resonances and the cut points for the broad group cross sections which are required by the GLEN code.

Normally cross sections for several temperatures are given on a PENDF. ETOGLEN will automatically process the cross sections for all temperatures given. These temperatures are initially chosen to span the range of interest and at intervals frequent enough to represent cross sections at a number of neighboring

temperatures. For the HTGR problem, for example, cross sections for 12 temperatures were required over a range from 300 to 3000 K. PENDF cross sections were generated at 0, 300, 950, and 3000 K, and Table B-I shows which of the PENDF values were used for each of the 12 temperatures. Table B-II describes the input specifications for ETOGLEN, and a listing of the code appears at the end of this appendix.

Sample results from ETOGLEN are shown in Table B-III and Figs. B-1 and B-2. These are for ^{233}U , MAT-1260; for this problem, an 86-point energy grid was previously optimized for the thermal resonances of ^{235}U . The graphical output from ETOGLEN (Figs. B-1 and B-2) demonstrates the accuracy with which the resonance structure is reproduced with the coarser grid, and the weighted averaging done in the code indicates the amount of error incurred in the multigroup cross sections by using the coarser grid (Table B-III). Also note in this table the small effect of temperature on average cross sections for this isotope and this energy group structure.

TABLE B-I
 PENDF TEMPERATURES USED FOR CROSS SECTIONS
 FOR END-OF-EQUILIBRIUM CYCLE (EOEC) HTGR CASES

HTGR-EOEC Case No.	Temperature (K)	PENDF Temperature Used for Thermal Cross Sections
1	300	300
2	500	300
3	600	300
4	800	950
5	1000	950
6	1200	950
7	1500	950
8	1700	950
9	2000	3000
10	2300	3000
11	2600	3000
12	3000	3000

TABLE B-II
INPUT SPECIFICATIONS FOR ETOGLEN

<u>Card No.</u>	<u>Format</u>	<u>Variable</u>	<u>Comment</u>
1	I11	NUMBIN	Number of energy mesh intervals over which equal lethargy intervals are specified.
2	6E11.4	BMIN(N), BMAX(N), DELU(N)	The minimum energy, the maximum energy, and the lethargy increment for each of the NUMBIN intervals.
3	I11	NPD	Number of additional energy points to be added.
4	6E11.4	ED(I)	NPD values of additional energies.
5	I11	NFGP	Number of cut points of intervals over which resonance integrals are to be computed (normally number of few groups, i.e., broad groups).
6	6E11.4	EC(I)	NFGP values of cut point energies. Note - if EC(I) are broad-group boundaries, they must also be specified in the ED list.
7	I11	NW	Number of energy-flux pairs given for the weighting function.
8	6E11.4	EWI(N), WI(N)	NW values for energy-flux pairs of specified weighting function.

NDF is name of the file containing pointwise data at several temperatures for the absorber being processed.

TABLE B-III

WEIGHTED AVERAGE FISSION
CROSS SECTIONS FOR ^{233}U

T = 300 K

<u>Upper Energy Boundary (eV)</u>	<u>Calculated Using Original PENDF Data</u>	<u>Calculated Using Data on Reduced Mesh</u>	<u>% Diff.</u>
0.04	566.6	567.5	0.2
0.10	324.9	328.5	1.1
0.414	201.6	205.3	1.8
2.38	232.9	234.6	0.7

T = 3000 K

0.04	566.7	567.5	0.2
0.10	324.9	328.6	1.1
0.414	202.1	205.8	1.8
2.38	232.2	233.4	0.5

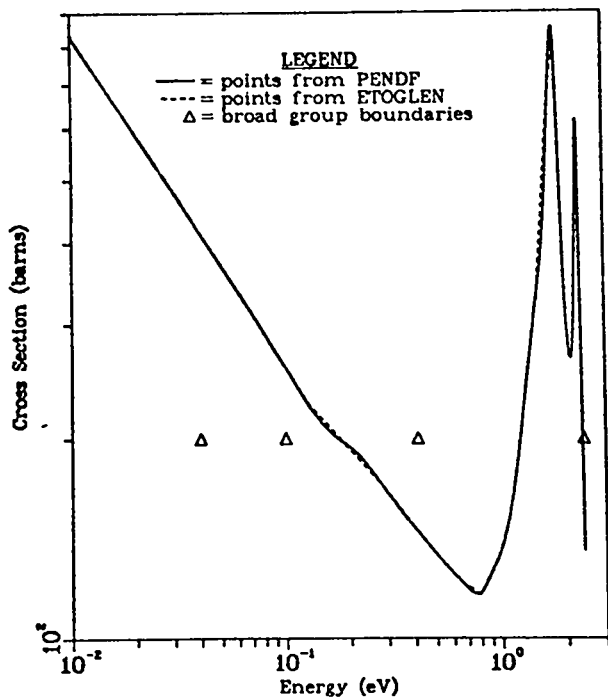


Fig. B-1.
 ^{233}U fission cross section at 300 K
from 0.01 to 2.38 eV.

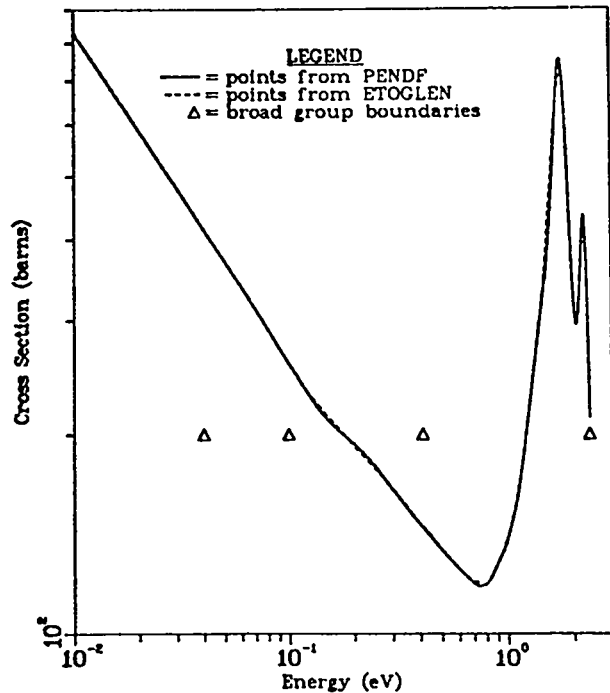


Fig. B-2.
 ^{233}U fission cross section at 3000 K
from 0.01 to 2.38 eV.

```

PROGRAM ETOGLEN(INP,OUT,PUN,FILM,FSET11,FSET12,FSET5=INP,      ETOGL  1
1          FSET6=OUT)                                          ETOGL  2
PROGRAM TO GET GLEN INPUT CROSS SECTIONS FROM PENDF TAPE,I.E. ETOGL  3
ENDF/B TO GLEN                                              ETOGL  4
-    --    ----                                             ETOGL  5
C                                                                ETOGL  6
C                                                                ETOGL  7
C                                                                ETOGL  8
C                                                                ETOGL  9
C                                                                ETOGL 10
C                                                                ETOGL 11
C                                                                ETOGL 12
C                                                                ETOGL 13
C                                                                ETOGL 14
C                                                                ETOGL 15
C                                                                ETOGL 16
C                                                                ETOGL 17
C                                                                ETOGL 18
C                                                                ETOGL 19
C                                                                ETOGL 20
C                                                                ETOGL 21
C                                                                ETOGL 22
C                                                                ETOGL 23
C                                                                ETOGL 24
C                                                                ETOGL 25
C                                                                ETOGL 26
C                                                                ETOGL 27
C                                                                ETOGL 28
C                                                                ETOGL 29
C                                                                ETOGL 30
C                                                                ETOGL 31
C                                                                ETOGL 32
C                                                                ETOGL 33
C                                                                ETOGL 34
C                                                                ETOGL 35
C                                                                ETOGL 36
C                                                                ETOGL 37
C                                                                ETOGL 38
C                                                                ETOGL 39
C                                                                ETOGL 40
C                                                                ETOGL 41
C                                                                ETOGL 42
C                                                                ETOGL 43
C                                                                ETOGL 44
C                                                                ETOGL 45
C                                                                ETOGL 46
C                                                                ETOGL 47
C                                                                ETOGL 48
C                                                                ETOGL 49
C                                                                ETOGL 50
C                                                                ETOGL 51
C                                                                ETOGL 52
C                                                                ETOGL 53
C                                                                ETOGL 54
C                                                                ETOGL 55
C                                                                ETOGL 56
C                                                                ETOGL 57
C                                                                ETOGL 58
C                                                                ETOGL 59
C                                                                ETOGL 60
C                                                                ETOGL 61
C                                                                ETOGL 62
DIMENSION XFISS(200),XCAP(200)
DIMENSION FC(200),ED(200),EM(200),INT(10),NPT(10),E(2000),S(2000),
1  HOL(10),SM(200),TITL(10),XLAB(10),YLAB(10),ER(2000),EMR(200),
2  SC(200),SR(2000),WI(200),W(2000),WM(200),EWI(200),NECT(10),
3  NEMCT(10),SBD(10),SMBD(10),BMIN(12),BMAX(12),DELU(12)

DIMENSION XP(500),YP(500)

E. IS BASIC CALCULATED MESH,ED ARE ADDITIONAL POINTS TO BE ADDED,
E.G. BROAD GROUP MESH BREAK POINTS IN GLEN,RESONANCE PEAKS,
VALLEYS,ETC.,AND EM IS COMBINED MESH.
FIRST READ BMIN,BMAX,DELU FOR EACH REGION (UP TO 8) FOR
CALCULATION OF BASIC E MESH -- DESCENDING ORDER.
EMIN IS LOWEST ENERGY BOUND, EMAX IS HIGHEST
EC ARE BROAD GROUP BREAK POINTS TO BE ADDED TO PENDF MESH
FOR INTEGRAL CHECK
EWI,WI ARE ENERGY,WEIGHT FUNCTION PAIRS FOR WEIGHTING
IF INTEGRAL CHECK

READ (5,30) NUMBIN
DO 5 N=1,NUMBIN
READ (5,10) BMIN(N),BMAX(N),DELU(N)
5 CONTINUE
10 FORMAT (6E11.4)
IF (NUMBIN.EQ.1) GO TO 7
NUMED=NUMBIN-1
DO 6 N=1,NUMED
IF (BMIN(N).NE.BMAX(N+1)) BMIN(N)=BMAX(N+1)
6 CONTINUE
7 NPC=1
WRITE (6,12) NUMBIN
12 FORMAT (14H1,17X,12,20H INPUT ENERGY GROUPS//2X,9HGROUP NO.,
1  2X,14HGROUP MAX (EV),2X,14HGROUP MIN (EV),2X,
2  14HLETHARGY WIDTH)
DO 8 N=1,NUMBIN
WRITE (6,13) N,BMAX(N),BMIN(N),DELU(N)
8 CONTINUE
13 FORMAT (4X,13,8X,E11.5,5X,E11.5,5X,E11.5)
DO 20 N=1,NUMBIN
U=0.
E(NPC)=BMAX(N)
11 U=U+DELU(N)
NPC=NPC+1
E(NPC)=BMAX(N)/EXP(U)
IF (E(NPC).LE.BMIN(N)) GO TO 20
GO TO 11
20 CONTINUE
E(NPC)=BMIN(N)
EMIN=BMIN(NUMBIN)
EMAX=BMAX(1)
DO 21 N=1,NPC
SR(N)=N
21 CONTINUE

READ ED MESH FROM CARDS.

```

	READ (5,30) NPD	ETOGL 63
30	FORMAT '(6I11)	ETOGL 64
	READ (5,10) (ED(I),I=1,NPD)	ETOGL 65
C		ETOGL 66
C	READ BROAD GROUP ENERGIES. READ FROM HI TO LO.	ETOGL 67
C		ETOGL 68
	READ (5,30) NFGP	ETOGL 69
	READ (5,10) (EC(I),I=1,NFGP)	ETOGL 70
C		ETOGL 71
C	READ IN WEIGHT FCNS. READ IN E-LO TO E-HI. MUST BE LOG-LOG INTERP.	ETOGL 72
C		ETOGL 73
	READ (5,30) NW	ETOGL 74
	READ (5,10) (EWI(N),WI(N),N=1,NW)	ETOGL 75
	TITL(1)=10W THERMAL WE	ETOGL 76
	TITL(2)=10W IGH T FUNCT	ETOGL 77
	TITL(3)=10W ION (GLEN)	ETOGL 78
	NWPLT=-NW	ETOGL 79
	XLAB(1)=10W ENERGY (EV)	ETOGL 80
	YLAB(1)=10W THERM FLUX	ETOGL 81
	CALL PLOJB (EWI,WI,NWPLT,-1,0,0,0,0,1,0,1,0,TITL,30,XLAB,10,YLAB,	ETOGL 82
	1 10)	ETOGL 83
C		ETOGL 84
C		ETOGL 85
C	COMBINE E AND ED TO FORM EM.	ETOGL 86
C		ETOGL 87
	CALL UNION (E,ED,NPC,NPD)	ETOGL 88
C		ETOGL 89
C	REORDER E TO GET EM	ETOGL 90
C		ETOGL 91
	DO 25 N=1,NPC	ETOGL 92
	NN1=NPC-N+1	ETOGL 93
	EM(N)=E(NN1)	ETOGL 94
25	CONTINUE	ETOGL 95
	JNPC=NPC-1	ETOGL 96
	NNTST=0	ETOGL 97
	DO 35 JI=1,JNPC	ETOGL 98
	DLTST=(EM(JI+1)-EM(JI))/EM(JI)*100.	ETOGL 99
	IF (DLTST.GT.1.0) GO TO 35	ETOGL100
	PRINT 36,JT,EM(JI),EM(JI+1),DLTST	ETOGL101
	NNTST=NNTST+1	ETOGL102
35	CONTINUE	ETOGL103
36	FORMAT (1H0,* DUPLICATE ENERGIES AT *14.* *1PE12.5,* AND *1PE12.5,	ETOGL104
1	* PCT DIFF = *F6.3)	ETOGL105
	IF (NNTST.FQ.0) PRINT 37	ETOGL106
37	FORMAT (1H0,* THERE ARE NO ENERGY PAIRS WITHIN ONE PERCENT,*)	ETOGL107
	NPM=NPC	ETOGL108
40	FORMAT (1H0,* I = *I3.* E = *1PE12.5)	ETOGL109
C		ETOGL110
C	PUNCH ENERGY MESH FOR GLEN.	ETOGL111
C		ETOGL112
	DO 45 N=1,NPM	ETOGL113
	NN1=NPM-N+1	ETOGL114
	E(N)=EM(NN1)	ETOGL115
45	CONTINUE	ETOGL116
	TITL(5)=10W ENERGY MES	ETOGL117
	TITL(6)=10W H FOR GLEN	ETOGL118
	PUNCH 210,(TITL(I),I=5,6)	ETOGL119
	PUNCH 150,(E(N),N=1,NPM)	ETOGL120
C		ETOGL121
C	THERE ARE SEVERAL TEMPERATURES ON TAPE. READ NOTEMP=NO. OF TEMPS.	ETOGL122
C		ETOGL123
	READ (5,30) NOTEMP,MAT1	ETOGL124
	DO 1000 NNT=1,NOTEMP	ETOGL125

C		ETOGL126
C	FOR EACH TEMP,GET XSEC FOR MT=2,MT=18,MT=102.	ETOGL127
C		ETOGL128
	DO 900 NMT=1,3	ETOGL129
	MF1=3	ETOGL130
	MT1=2	ETOGL131
	IF (NMT.EQ.2) MT1=18	ETOGL132
	IF (NMT.EQ.3) MT1=102	ETOGL133
	IF (NMT.EQ.3.AND.MAT1.EQ.1155) MT1=107	ETOGL134
50	READ (11,60) (HOL(I),I=1,7),MAT,MF,MT,NSEQ	ETOGL135
60	FORMAT (A6,10,A6,I4,I2,I3,I5)	ETOGL136
	IF (MAT.LT.0) GO TO 2000	ETOGL137
	IF (MAT.LT.MAT1) GO TO 50	ETOGL138
	IF (MAT.GT.MAT1) GO TO 2000	ETOGL139
	IF (MF.NE.3) GO TO 50	ETOGL140
	IF (MT.NE.MT1) GO TO 50	ETOGL141
	WRITE (6,2020) MAT,MAT1,MF,MT,MT1,NMT	ETOGL142
2020	FORMAT (1H,*MAT=*I4,3X,*MAT1=*I4,3X,*MF=*I2,3X,*MT=*I3,3X,*MT1=*,	ETOGL143
1	I3,3X,*NMT=*I2)	ETOGL144
	READ (11,70) C1,C2,N1,N2,NR,NP	ETOGL145
70	FORMAT(1P2E11.4,4I11)	ETOGL146
	TEMDS=C1	ETOGL147
	READ (11,30) (NPT(I),INT(I),I=1,NR)	ETOGL148
C		ETOGL149
C	ASSUME THERMAL RANGE IS WITHIN FIRST 2000 PTS ON TAPE.	ETOGL150
C		ETOGL151
	NPTH=NP	ETOGL152
	IF (NP.GT.2000) NPTH=2000	ETOGL153
	READ (11,10) (E(I),S(I),I=1,NPTH)	ETOGL154
80	READ (11,60) (HOL(I),I=1,7),MAT,MF,MT,NSEQ	ETOGL155
	IF (MT.NE.0) GO TO 80	ETOGL156
	IF (E(NPTH).GT.EMAX) GO TO 82	ETOGL157
	WRITE (6,81) NPTH,E(NPTH)	ETOGL158
81	FORMAT (1H1,* EMAX NOT WITHIN *I4,* PTS. LAST ENERGY = *1PE12.5)	ETOGL159
	STOP	ETOGL160
82	CONTINUE	ETOGL161
C		ETOGL162
C	GET XSEC,SM,CORRESPONDING TO EM.	ETOGL163
C		ETOGL164
	DO 100 I=1,NPM	ETOGL165
	ILO=LOCT(E,EM(I),NPTH)	ETOGL166
	IF (ILO.EQ.-1) CALL ERROR(100)	ETOGL167
	IHI=ILO+1	ETOGL168
	DO 85 J=1,NR	ETOGL169
	IF (IHI.LE.NPT(J)) GO TO 90	ETOGL170
85	CONTINUE	ETOGL171
	CALL ERROR (200)	ETOGL172
90	CALL TERP1 (E(ILO),S(ILO),E(IHI),S(IHI),EM(I),CSEC,INT(J))	ETOGL173
	SM(I)=CSEC	ETOGL174
100	CONTINUE	ETOGL175
110	FORMAT (1H0,* I=*I6,* EM = *1PE11.4,* SM = *1PE11.4)	ETOGL176
C		ETOGL177
C	CHECK INTEGRALS AND MAKE COMPARISON PLOTS.	ETOGL178
C		ETOGL179
C	CUT OFF MESH POINTS ABOVE EMAX	ETOGL180
	KTHRM=0	ETOGL181
	DO 120 N=1,NPTH	ETOGL182
	IF (E(N).GT.EMAX) GO TO 130	ETOGL183
	KTHRM=KTHRM+1	ETOGL184
120	CONTINUE	ETOGL185
130	CONTINUE	ETOGL186
	TITL(1)=10HETOGLEN VS	ETOGL187
	TITL(2)=10H PENDF PTS	ETOGL188

	TITL(3)=10H. ELASTIC	ETOGL189
	TITL(4)=10HCROSS SECT	ETOGL190
	IF (NMT.FQ.2) TITL(3)=10H. FISSION	ETOGL191
	IF (NMT.FQ.3) TITL(3)=10H. CAPTURE	ETOGL192
	NPTR=KTHRM	ETOGL193
C		ETOGL194
C	REVERSE E-MESH AND ADD BROAD GROUP CUT POINTS	ETOGL195
C		ETOGL196
	DO 300 N=1,NPTR	ETOGL197
	N1=NPTR-N+1	ETOGL198
	ER(N)=E(N1)	ETOGL199
300	CONTINUE	ETOGL200
	CALL UNION(ER,EC,NPTR,NFGP)	ETOGL201
	DO 320 N=1,NPTR	ETOGL202
	ILO=LOCT(E,ER(N),NPTR)	ETOGL203
	IF (ILO.EQ.-1) CALL ERROR(200)	ETOGL204
	IHI=ILO+1	ETOGL205
	CALL TERP1 (E(ILO),S(ILO),E(IHI),S(IHI),ER(N),CSEC,2)	ETOGL206
	SR(N)=CSEC	ETOGL207
320	CONTINUE	ETOGL208
C		ETOGL209
C	CUT OFF MESH POINTS BELOW EMIN	ETOGL210
C		ETOGL211
321	DO 326 N=1,NPTR	ETOGL212
	IF (ER(N).LT.EMIN) GO TO 327	ETOGL213
	NSTOR=N	ETOGL214
326	CONTINUE	ETOGL215
327	CONTINUE	ETOGL216
	IF (ER(NSTOR).EQ.EMIN) N=NSTOR+1	ETOGL217
	NPTR=N-1	ETOGL218
	KTHRM=NPTR	ETOGL219
C		ETOGL220
C	REORDER LOW TO HIGH	ETOGL221
C		ETOGL222
	DO 330 N=1,NPTR	ETOGL223
	N1=NPTR-N+1	ETOGL224
	E(N)=ER(N1)	ETOGL225
	S(N)=SR(N1)	ETOGL226
330	CONTINUE	ETOGL227
C	PUT WT FCN ON E AND EM MESHES.	ETOGL228
C		ETOGL229
	DO 340 N=1,NPTR	ETOGL230
	ILO=LOCT(EWI,E(N),NW)	ETOGL231
	IF (ILO.EQ.-1) CALL ERROR (300)	ETOGL232
	IHI=ILO+1	ETOGL233
	CALL TERP1 (EWI(ILO),WI(ILO),EWI(IHI),WI(IHI),E(N),WSS,5)	ETOGL234
	W(N)=WSS	ETOGL235
340	CONTINUE	ETOGL236
	DO 350 N=1,NPM	ETOGL237
	ILO=LOCT(EWI,EM(N),NW)	ETOGL238
	IF (ILO.EQ.-1) CALL ERROR(400)	ETOGL239
	IHI=ILO+1	ETOGL240
	CALL TERP1 (EWI(ILO),WI(ILO),EWI(IHI),WI(IHI),EM(N),WSS,5)	ETOGL241
	WM(N)=WSS	ETOGL242
350	CONTINUE	ETOGL243
C		ETOGL244
C	REVERSE EC MESH	ETOGL245
C		ETOGL246
	DO 345 K=1,NFGP	ETOGL247
	SC(K)=EC(K)	ETOGL248
345	CONTINUE	ETOGL249
	DO 346 K=1,NFGP	ETOGL250
	K1=NFGP-K+1	ETOGL251

	EC(K)=SC(K1)	ETOGL252
346	CONTINUE	ETOGL253
C		ETOGL254
C	FIND BROAD GROUP CUT POINTS IN E AND EM MESHES.	ETOGL255
C		ETOGL256
	NN1=0	ETOGL257
	DO 370 I=1,NFGP	ETOGL258
	DO 360 N=1,NPM	ETOGL259
	IF ((EM(N)-EC(I)).NE.0,0) GO TO 360	ETOGL260
	NN1=NN1+1	ETOGL261
	NEMCT(NN1)=N	ETOGL262
	GO TO 370	ETOGL263
360	CONTINUE	ETOGL264
370	CONTINUE	ETOGL265
	WRITE (6,3A0) (NEMCT(N),N=1,NFGP)	ETOGL266
3A0	FORMAT (1H0,* BROAD GROUP CUT POINTS */5X,10I6)	ETOGL267
	NN1=0	ETOGL268
	DO 400 I=1,NFGP	ETOGL269
	DO 390 N=1,NPTR	ETOGL270
	IF (E(N).NE.EC(I)) GO TO 390	ETOGL271
	NN1=NN1+1	ETOGL272
	NECT(NN1)=N	ETOGL273
	GO TO 400	ETOGL274
390	CONTINUE	ETOGL275
400	CONTINUE	ETOGL276
	WRITE (6,3A0) (NECT(I),I=1,NFGP)	ETOGL277
	WRITE (6,405) (TITL(I),I=3,4),NNT	ETOGL278
405	FORMAT (1H1,* GLEN POINTS FOR *2A10,* . TEMP NUMBER *,I2,	ETOGL279
	1 /** PT. NO,*4X,*ENERGY*8X,	ETOGL280
	2 *CROSS SECTION*4X,*WEIGHT FUNCTION*)	ETOGL281
407	FORMAT (1H1,* PENDF POINTS FOR *2A10,* . TEMP NUMBER *,I2,	ETOGL282
	1 /** PT. NO,*4X,*ENERGY*8X,	ETOGL283
	2 *CROSS SECTION*4X,*WEIGHT FUNCTION*)	ETOGL284
	WRITE (6,406) (N,EM(N),SM(N),WM(N),N=1,NPM)	ETOGL285
	WRITE (6,407) (TITL(I),I=3,4),NNT	ETOGL286
406	FORMAT (16.1P3E18.5)	ETOGL287
	WRITE (6,406) (N,E(N),S(N),W(N),N=1,NPTR)	ETOGL288
C		ETOGL289
C	GET BROAD GROUP XSEC FOR BOTH PENDF AND GLEN DATA.	ETOGL290
C		ETOGL291
	NBG=NFGP-1	ETOGL292
	DO 430 N=1,NBG	ETOGL293
	TOP=0.	ETOGL294
	DEM=0.	ETOGL295
	NE1=NECT(N)	ETOGL296
	NE2=NECT(N+1)-1	ETOGL297
	NEM1=NEMCT(N)	ETOGL298
	NEM2=NEMCT(N+1)-1	ETOGL299
	DO 410 J=NE1,NE2	ETOGL300
	TOP=TOP+(E(J+1)-E(J))*(W(J+1)*S(J+1)+W(J)*S(J))/2.	ETOGL301
	DEM=DEM+((E(J+1)-E(J))*(W(J+1)+W(J))/2,0)	ETOGL302
410	CONTINUE	ETOGL303
	SBD(N)=TOP/DEM	ETOGL304
	TOP=0.	ETOGL305
	DEM=0.	ETOGL306
	DO 420 J=NEM1,NEM2	ETOGL307
	TOP=TOP+(EM(J+1)-EM(J))*(WM(J+1)*SM(J+1)+WM(J)*SM(J))/2.	ETOGL308
	DEM=DEM+((EM(J+1)-EM(J))*(WM(J+1)+WM(J))/2,0)	ETOGL309
420	CONTINUE	ETOGL310
	SMBD(N)=TOP/DEM	ETOGL311
430	CONTINUE	ETOGL312
	WRITE (6,440) (TITL(I),I=3,4),NNT	ETOGL313
440	FORMAT (1H1,20X,2A10,* . TEMP NUMBER *I2,	ETOGL314

```

1      /* BROAD GROUP ENERGY XSEC FROM PENDF XSEC FROM GLEN DATA*/
WRITE (6,450) (EC(I),SBD(I),SMRD(I),I=1,NB8)
450  FORMAT (1P3E18,5)
      NI=1
      XLAB(1)=10H ENERGY IN
      XLAB(2)=10H EV. UNITS
      YLAB(1)=10H CROSS SEC
      YLAB(2)=10HTION (BNS)
      NPLOT=0
      DO 455 N=1,KTHRM
      IF (E(N).LT.0.01) GO TO 455
      NPLOT=NPLOT+1
      XP(NPLOT)=E(N)
      YP(NPLOT)=S(N)
455  CONTINUE
      NPLOT=-NPLOT
      NI=-1
      CALL PLOTM (XP,YP,NPLOT,NI,0,39,0.,1.,1.,TITL,40,XLAB,20,
1      YLAB,20)
      NPLT=NPM
      NPLOT=0
      DO 460 N=1,NPLT
      IF (EM(N).LT.0.01) GO TO 460
      NPLOT=NPLOT+1
      XP(NPLOT)=EM(N)
      YP(NPLOT)=SM(N)
460  CONTINUE
      NPLOT=-NPLOT
      NI=-1
      CALL PLOTM (XP,YP,NPLOT,NI,0,-37,0.,1.,1.,TITL,40,XLAB,20,
1      YLAB,20)

C
C      REORDER FOR GLEN PUNCH
C
      DO 140 N=1,NPM
      NN1=NPM-N+1
      E(N)=EM(NN1)
      S(N)=SM(NN1)
140  CONTINUE

C
C      PUNCH FOR GLEN
C
      TITL(1)=10H ELASTIC
      TITL(2)=10HCROSS SECT
      TITL(3)=10HION FOR MA
      TITL(4)=10HTERIAL *
      IF (NMT.NE.1) GO TO 141
      PUNCH 200,(TITL(I),I=1,4),MAT1,TEMDS
      PUNCH 150,(S(N),N=1,NPM)
150  FORMAT (1P4E20,12)
200  FORMAT (4A10,I4,* TEMP=*1PE11,4,*DEG K*)
210  FORMAT (4A10)
141  CONTINUE
      IF (NMT.NE.2) GO TO 143
      DO 142 N=1,NPM
      XFISS(N)=S(N)
142  CONTINUE
143  CONTINUE
      IF (NMT.NE.3) GO TO 147
      DO 144 N=1,NPM
      XCAP(N)=S(N)
144  CONTINUE
      DO 145 N=1,NPM

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ETOGL315
ETOGL316
ETOGL317
ETOGL318
ETOGL319
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ETOGL369
ETOGL370
ETOGL371
ETOGL372
ETOGL373
ETOGL374
ETOGL375
ETOGL376
ETOGL377

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	S(N)=XFISS(N)*XCAP(N)	ETOGL378
145	CONTINUE	ETOGL379
	TITL(1)=10H ABSORPT	ETOGL380
	PUNCH 200,(TITL(I),I=1,4),MAT1,TEMDS	ETOGL381
	PUNCH 150,(S(N),N=1,NPM)	ETOGL382
	XNU=1.0	ETOGL383
	IF (MAT1.EQ.1157) XNU=2.4188	ETOGL384
	IF (MAT1.EQ.1260) XNU=2.498	ETOGL385
	DO 146 N=1,NPM	ETOGL386
	S(N)=XNU*XFISS(N)	ETOGL387
146	CONTINUE	ETOGL388
	TITL(1)=10H NUFISSN	ETOGL389
	PUNCH 200,(TITL(I),I=1,4),MAT1,TEMDS	ETOGL390
	PUNCH 150,(S(N),N=1,NPM)	ETOGL391
147	CONTINUE	ETOGL392
C		ETOGL393
C	REVERSE EC MESH FOR NEXT PASS.	ETOGL394
C		ETOGL395
	DO 211 K=1,NFGP	ETOGL396
	SC(K)=EC(K)	ETOGL397
211	CONTINUE	ETOGL398
	DO 212 K=1,NFGP	ETOGL399
	K1=NFGP-K+1	ETOGL400
	EC(K)=SC(K1)	ETOGL401
212	CONTINUE	ETOGL402
900	CONTINUE	ETOGL403
1000	CONTINUE	ETOGL404
	STOP	ETOGL405
2000	WRITE (6,2010) MAT1,MAT	ETOGL406
2010	FORMAT (1H1,* SORRY,MAT = *I4,* NOT ON TAPE. LAST MAT = *I4)	ETOGL407
	STOP	ETOGL408
	END	ETOGL409

	FUNCTION LOCT (E,EK,N)	LOCT 1
C	BRACKETS EK IN E SO THAT EK.GE.E(LOCT) AND EK.LT.E(LOCT+1)	LOCT 2
C	IF EK CANNOT BE BRACKETED, LOCT=-1	LOCT 3
	DIMENSION E(1)	LOCT 4
C	RETURN LOCT=-1 IF ARRAY HAS ONLY ONE PT. (AS FOR NPTS IN XSEC).	LOCT 5
50	FORMAT (1H0,1P8E15.5)	LOCT 6
	IF (N.LE.1) GO TO 10	LOCT 7
	M=N-1	LOCT 8
	DO 5 I=1,M	LOCT 9
5	IF ((EK.GE.E(I)).AND.(EK.LT.E(I+1))) GO TO 15	LOCT 10
	IF ((E(N-1).EQ.E(N)).AND.(EK.EQ.E(N))) GO TO 10	LOCT 11
	IF (E(N).NE.EK) GO TO 10	LOCT 12
	LOCT=M	LOCT 13
	RETURN	LOCT 14
10	LOCT=-1	LOCT 15
	PRINT 50,(E(I),I=1,N),EK	LOCT 16
	RETURN	LOCT 17
15	LOCT=I	LOCT 18
	RETURN	LOCT 19
	END	LOCT 20

C	SUBROUTINE TERP1 (X1,Y1,X2,Y2,X,Y,I)	TERP1 1
C	====INTERPOLATE ONE PT.====	TERP1 2
C	(X1,Y1) AND (X2,Y2) ARE END PTS. OF THE LINE	TERP1 3
C	(X,Y) IS INTERPOLATED POINT	TERP1 4
C	I=INTERPOLATION CODE	TERP1 5
C	NOTE - IF A NEGATIVE OR ZERO ARGUMENT OF A LOG IS DETECTED, THE	TERP1 6
C	INTERPOLATION IS AUTOMATICALLY CHANGED FROM LOG TO LINEAR.	TERP1 7
C	ERROR STOPS - 301 (X1=X2,DISCONTINUITY)	TERP1 8
C	302 (INTERPOLATION CODE IS OUT OF RANGE)	TERP1 9
C	303 (ZERO OR NEGATIVE ARGUMENT FOR INTERPOLATED PT.)	TERP1 10
	5 XA=X1	TERP1 11
	YA=Y1	TERP1 12
	XB=X2	TERP1 13
	YB=Y2	TERP1 14
	XP=X	TERP1 15
	II=I	TERP1 16
	IF (XA.EQ.XB) CALL ERROR (301)	TERP1 17
	IF (II) 10,10,15	TERP1 18
	10 CALL ERROR (302)	TERP1 19
	15 IF (II-5) 20,20,10	TERP1 20
	20 GO TO (25,30,35,60,75), II	TERP1 21
	25 YP=YA	TERP1 22
	IF (XP.EQ.XB) YP=YB	TERP1 23
	GO TO 105	TERP1 24
	30 YP=YA+(XP-XA)*(YB-YA)/(XB-XA)	TERP1 25
	GO TO 105	TERP1 26
	35 IF (XA) 30,30,40	TERP1 27
	40 IF (XB) 30,30,45	TERP1 28
	45 IF (XP) 50,50,55	TERP1 29
	50 CALL ERROR (303)	TERP1 30
	55 YP=YA+ALOG(XP/XA)*(YB-YA)/ALOG(XB/XA)	TERP1 31
	GO TO 105	TERP1 32
	60 IF (YA) 30,30,65	TERP1 33
	65 IF (YB) 30,30,70	TERP1 34
	70 YP=YA*EXP((XP-XA)*ALOG(YB/YA)/(XB-XA))	TERP1 35
	GO TO 105	TERP1 36
	75 IF (YA) 35,35,80	TERP1 37
	80 IF (YB) 35,35,85	TERP1 38
	85 IF (XA) 70,70,90	TERP1 39
	90 IF (XB) 70,70,95	TERP1 40
	95 IF (XP) 80,50,100	TERP1 41
	100 YP=YA*EXP(ALOG(XP/XA)*ALOG(YB/YA)/ALOG(XB/XA))	TERP1 42
	105 Y=YP	TERP1 43
	RETURN	TERP1 44
	END	TERP1 45

	SUBROUTINE ERROR (N)	ERROR 1
	IOS=9	ERROR 2
	5 PRINT 10,N	ERROR 3
	WRITE(99,10)	ERROR 4
	10 FORMAT (11H ERROR STOP,I6)	ERROR 5
	END	ERROR 6

	SUBROUTINE UNION (XU,X,NPU,NP)	UNION 1
C		UNION 2
C	FUNCTION OF SUBROUTINE	UNION 3
C	UNION1 COMPUTES THE UNION OF INDEPENDENT VARIABLE SETS X(IP),IP=1,UNION	UNION 4
C	XU(IP1),IP1=1,NPU, AND PLACES THE UNION INTO XU(IP2),IP2=1,NPUNION	UNION 5
C	STORAGE	UNION 6
	DIMENSION XU(2000),KU(2000),X(200)	UNION 7
C		UNION 8
C	ADD A SET X TO AN EXISTING UNION SET XU	UNION 9
	DO 106 IP=1,NPU	UNION 10
	KU(IP)=0	UNION 11
106	CONTINUE	UNION 12
	DO 103 IP=1,NP	UNION 13
	IF (X(IP).LT.XU(NPU)) GO TO 120	UNION 14
	IF (X(IP).GT.XU(1)) GO TO 130	UNION 15
	DO 104 IP1=1,NPU	UNION 16
	IF (X(IP).EQ.XU(IP1)) GO TO 140	UNION 17
	IF (IP1.EQ.NPU) GO TO 105	UNION 18
	IF (X(IP).LT.XU(IP1).AND.X(IP).GT.XU(IP1+1)) GO TO 150	UNION 19
105	CONTINUE	UNION 20
104	CONTINUE	UNION 21
C		UNION 22
C	HERE NPU IS INCREMENTED BY ONE AND A POINT IS ADDED TO THE LEFT	UNION 23
120	NPU=NPU+1	UNION 24
	XU(NPU)=X(IP)	UNION 25
	KU(NPU)=1	UNION 26
121	CONTINUE	UNION 27
	GO TO 103	UNION 28
C		UNION 29
C	HERE CONTROLS ARE SET TO ADD A POINT ON THE RIGHT	UNION 30
130	KONREL=1	UNION 31
	NPMOV=NPU	UNION 32
	GO TO 170	UNION 33
C		UNION 34
C	HERE NPU IS NOT INCREMENTED BY ONE .	UNION 35
140	CONTINUE	UNION 36
	KU(IP1)=1	UNION 37
	GO TO 103	UNION 38
C		UNION 39
C	HERE NPU IS INCREMENTED BY ONE AND CONTROLS ARE SET TO ADD A POINT	UNION 40
C	BETWEEN POINTS IP1 AND IP1+1	UNION 41
150	KONREL=2	UNION 42
	NPMOV=NPU-IP1	UNION 43
	GO TO 170	UNION 44
C		UNION 45
C	HERE WE INCREMENT NPU BY ONE AND MOVE THE LEFT-MOST NPMOV POINTS I	UNION 46
C	SET ONE POSITION TO THE LEFT	UNION 47
170	NPU=NPU+1	UNION 48
	DO 171 IP2=1,NPMOV	UNION 49
	XU(NPU-IP2+1)=XU(NPU-IP2)	UNION 50
	KU(NPU-IP2+1)=KU(NPU-IP2)	UNION 51
171	CONTINUE	UNION 52
C		UNION 53
C	HERE A NEW POINT IS ADDED	UNION 54
	NPADD=NPU-NPMOV	UNION 55
	XU(NPADD)=X(IP)	UNION 56
	KU(NPADD)=1	UNION 57
172	CONTINUE	UNION 58
103	CONTINUE	UNION 59
102	RETURN	UNION 60
	END	UNION 61

APPENDIX C

MERGFAT

A code to merge fast and thermal cross section sets.

Multigroup cross-section data sets for energy groups above the thermal boundary energy (2.58 eV for the HTGR) are generated by the IDX code, whereas data for those groups below this energy are generated by the GLEN code (see Fig. 3). Usually there is at least one overlapping group. MERGFAT is a small code, the purpose of which is to combine these two sets into a single set properly formatted for input to the Los Alamos S_n codes.

In addition to the files containing the fast and thermal data, the input consists of designations of groups to be merged, designations of materials to be read from the files, the final energy boundaries (which are used in the computation of the velocities needed by the S_n codes), and the final groupwise values of the fraction of the fissions in each group (χ), also needed in the S_n calculations. Table C-I describes the input needed for MERGFAT, and a listing of the code is given at the end of this appendix.

TABLE C-I

INPUT SPECIFICATIONS FOR MERGFAT

<u>Card No.</u>	<u>Format</u>	<u>Variable</u>	<u>Comment</u>
1	9A8	HLT(I)	Should read "LASTDECK" at the last set of input cards.
2	12I6	LENG	Table length of final output, including upscatter, self-scatter, and down-scatter.
		NDELU	Obsolete.
		LTL	Obsolete.
		LNGUP	Length of up-scatter table.
3	8A10	TITLE(I)	Title card.
4	3I12	NOBG	Total number of final groups.
		NOI	Number of materials for which cross sections are to be prepared.
		IOPT	Obsolete.

TABLE C-1 (cont)

5	6E12.5	GPEN(I)	NOBG values of lower group bounds in eV.
6	6E12.5	XI(I)	NOBG values of X.
7	12I6	NDKS	Number of sets to be merged (2).
		KG1(N),KG2(N)	NDKS values for first and last group in the sets being merged that are to be included in the final set.
8	A6,2X,A10,A6	NUCLE	Nuclide identifier assigned in 1DX input (see Ref. 14).
		MODER	Moderator, absorber identifier. Use word "MODERAT" for moderator and "ISOTOPE" for absorber.
		MATID	ENDF/B MAT number.

The "fast" data file output by the 1DX code is designated as NTPF in MERGFAT, and the "thermal" data file is designated as NTPT. Card 8 is repeated NOI times for the number of materials to be processed in one run.

LASL Identification No. LP-0757

```

PROGRAM MERGFAT (INP,OUT,PUN,FSET6=OUT,FSET7=PUN,FSET8=INP,
1 FSET9,FSET10) MERGF 1
C MERGFAT IS A VERSION OF JUMRLFAT THAT ACCEPTS FAST XSEC AS MERGF 2
C OUTPUT BY THE HONEDX VERSION OF THE 1DX CODE. MERGF 3
C REVISIONS MADE AT LASL BY LABAUVE,NOV75. MERGF 4
C MERGF 5
C MERGF 6
DIMENSION TOTP0(70),TOTP1(70),TOTRA(70) MERGF 7
DIMENSION NUCID(20),N2N(20),SK2KF(70,70),DIAGSM(70),CHKSM(70),
1 DIFF(70),CAPA(70),FISA(70),STR(70),AVNU(70),SINTRA(70,70), MERGF 8
2 SCAP(70),SS(70),SABS(70),HOL(70),KK(70),SSN2N(70,70), MERGF 9
3 XNUSYG(70),A(2000),P0(70,70),P1(70,70),P2(70,70),P3(70,70), MERGF 10
4 XS(70),TITLE(12),TOTN2N(70),SNP(70),SND(70),SNT(70),SNHE3(70), MERGF 11
5 SNA(70),SN2A(70) MERGF 12
DIMENSION XI(50),GPEN(50),TTL(12),VEL(50),FACAP(50),FAFIS(50), MERGF 13
1 ADEN(20),TOTIN(70) MERGF 14
DIMENSION KG1(20),KG2(20) MERGF 15
DIMENSION V(50),HLT(10) MERGF 16
NTPF=9 MERGF 17
NTPT=10 MERGF 18
5002 READ(8,5001)(HLT(I),I=1,9) MERGF 19
5001 FORMAT(9A8) MERGF 20
PRINT 600 MERGF 21
600 FORMAT(1H12OX10HINPUT DATA ///) MERGF 22
READ(8,5) LENG,NDELU,LTL,LANGUP MERGF 23
PRINT 601,LENG,NDELU,LTL,LANGUP MERGF 24
MERGF 25

```

601	FORMAT (1H010X7HLENG = ,I6,10H, NDELU = ,I6,8H, LTL = ,I6,	MERGF	26
	110H, LNGUP = ,I6)	MERGF	27
	READ (8,70) (TITLE(I),I=1,7)	MERGF	28
	PRINT 602,(TITLE(I),I=1,7)	MERGF	29
602	FORMAT (1H05X,8A10)	MERGF	30
	READ (8,71) NOBG,NOI,IOPT	MERGF	31
	PRINT 603,NOBG,NOI,IOPT	MERGF	32
603	FORMAT (1H010X7HNOBG = ,I12,8H, NOI = ,I12,9H, IOPT = ,I12)	MERGF	33
	NMAT=NOI	MERGF	34
71	FORMAT (3I12)	MERGF	35
	KGROPS=NOBG	MERGF	36
70	FORMAT (8A10)	MERGF	37
	PUNCH 1000,(TITLE(I),I=1,12)	MERGF	38
	READ (8,3) (GPN(I),I=1,KGROPS)	MERGF	39
	PRINT 604	MERGF	40
604	FORMAT (4H0 13X7HGPEN(I))	MERGF	41
	PRINT 605,(I,GPEN(I),I=1,KGROPS)	MERGF	42
605	FORMAT (13,1PE12.5)	MERGF	43
	AVEN1=(GPEN(KGROPS-1)*GPEN(KGROPS))/2.0	MERGF	44
	AVEN2=(GPEN(KGROPS))/2.0	MERGF	45
	FAXT=0.	MERGF	46
	READ (8,3) (XI(I),I=1,KGROPS)	MERGF	47
	PRINT 606	MERGF	48
606	FORMAT (4H0 16X5HXI(I))	MERGF	49
	PRINT 605,(I,XI(I),I=1,KGROPS)	MERGF	50
	V(1)=1.0F+07	MERGF	51
	DO 4000 K=1,KGROPS	MERGF	52
4000	V(K+1)=GPFN(K)	MERGF	53
	DO 72 K=1,KGROPS	MERGF	54
72	VEL(K)=0.007*(V(K)**0.5+V(K+1)**0.5)	MERGF	55
C	*** NOTE *** THESE ARE JUST AVERAGE BROAD GROUP VELOCITIES.	MERGF	56
	LBL1=6H SORS	MERGF	57
	LBL2=6H VELS	MERGF	58
	I1=1	MERGF	59
	M1=1	MERGF	60
73	I2=I1+5	MERGF	61
	IF (I2.LE.KGROPS) GO TO 74	MERGF	62
	PUNCH 75,(XI(I),I=I1,KGROPS)	MERGF	63
75	FORMAT (6F12.6)	MERGF	64
	GO TO 76	MERGF	65
74	PUNCH 87,(XI(I),I=I1,I2),LRL1,M1	MERGF	66
	M1=M1+1	MERGF	67
	I1=I2+1	MERGF	68
	GO TO 73	MERGF	69
76	CONTINUE	MERGF	70
	I1=1	MERGF	71
	M1=1	MERGF	72
77	I2=I1+5	MERGF	73
	IF (I2.LE.KGROPS) GO TO 78	MERGF	74
	PUNCH 79,(VEL(I),I=I1,KGROPS)	MERGF	75
79	FORMAT (1P6E12.2)	MERGF	76
	GO TO 1180	MERGF	77
78	PUNCH 81,(VEL(I),I=I1,I2),LBL2,M1	MERGF	78
	M1=M1+1	MERGF	79
	I1=I2+1	MERGF	80
	GO TO 77	MERGF	81
1180	CONTINUE	MERGF	82
87	FORMAT (6F12.6,A6,I2)	MERGF	83
81	FORMAT (1P6E12.2,A6,I2)	MERGF	84
5	FORMAT (12I6)	MERGF	85
	MATNO=1	MERGF	86
	READ (8,5) NDKS,(KG1(N),KG2(N),N=1,NDKS)	MERGF	87
	PRINT 607,NDKS	MERGF	88

607	FORMAT (1H010X7HNDKS = ,I6)	MERGF 89
	PRINT 608	MERGF 90
608	FORMAT (4H0 N3X6HKG1(N)7X6HKG2(N))	MERGF 91
	PRINT 609, (N,KG1(N),KG2(N),N=1,NOKS)	MERGF 92
609	FORMAT (I3.2I6)	MERGF 93
	DO 999 M=1,NMAT	MERGF 94
	REWIND NTPF \$ REWIND NTPT	MERGF 95
	READ (8,4100) NUCLE,MODER,MATID	MERGF 96
4100	FORMAT (A6,2X,A10,A6)	MERGF 97
	DO 500 N=1,NOKS	MERGF 98
	K1=KG1(N)	MERGF 99
	NOBG=KG2(N)	MERGF100
	LTABL=LENG+6	MERGF101
	NOBG2 = NOBG + 1	MERGF102
	FAFIS(N)=1.0	MERGF103
2102	FACAP(N)=1.	MERGF104
	R5 FORMAT (15H G=03,3H L=03,13H NUCLE NO.=05,1H 3A10)	MERGF105
4110	IF (N.EQ.1) READ (NTPF,4115) NUCID(M),FID,(TTL(KK),KK=1,11)	MERGF106
4115	FORMAT (13A6)	MERGF107
	IF (NUCID(M).NE.NUCLE) GO TO 4110	MERGF108
1085	FORMAT (A6,F6.2,11A6)	MERGF109
	DO 3070 K=K1,NOBG	MERGF110
	DO 3070 KF=K1,NOBG	MERGF111
	P0(K,KF)=0.0	MERGF112
	SINTRA(K,KF)=0.0	MERGF113
	SSN2N(K,KF)=0.0	MERGF114
3070	CONTINUE	MERGF115
	IF (N.EQ.1) GO TO 3000	MERGF116
	GNU=AVNU(K1)	MERGF117
	MTBL=2*LENGUP+5	MERGF118
	MGPS=NOBG-K1+1	MERGF119
	MVLG=MTBL*MGPS	MERGF120
	DO 3010 L=1,2	MERGF121
4120	READ (NTPT,4130) MODES,MATCH	MERGF122
4130	FORMAT (22X,A10,15X,A6)	MERGF123
	IF (MODES.EQ.MODER.AND.MATCH.EQ.MATID) GO TO 4135	MERGF124
	GO TO 4120	MERGF125
115	.03	MERGF126
4135	CONTINUE	MERGF127
	IF (L.EQ.1) MADD=0	MERGF128
	IF (L.EQ.2) MADD=MVLG	MERGF129
	READ (NTPT,3) (A(M*MADD),M=1,MVLG)	MERGF130
3010	CONTINUE	MERGF131
	DO 3020 MG=1,MGPS	MERGF132
	DO 3030 MGM=1,MGPS	MERGF133
	KK1=K1+MG-1	MERGF134
	KK2=K1+MGM-1	MERGF135
	LCK=9-MG+(MGM-1)*(MTBL+1)	MERGF136
	P0(KK1,KK2)=A(LCK)	MERGF137
	P1(MG,MGM)=A(LCK+MVLG)	MERGF138
3030	CONTINUE	MERGF139
	PRINT 3040,KK1	MERGF140
3040	FORMAT (1H0,* SCATT, MATRIX FOR GP *I3)	MERGF141
	PRINT 3050, (P0(KK1,KK2),KK2=K1,NOBG)	MERGF142
	PRINT 3050, (P1(MG,MGM),MGM=1,MGPS)	MERGF143
	LCK=(MG-1)*MTBL+1	MERGF144
	SABS(KK1)=A(LCK)	MERGF145
	XNUSIG(KK1)=A(LCK+1)	MERGF146
	STR(KK1)=A(LCK+2)	MERGF147
	AVNU(KK1)=GNU	MERGF148
	IF (GNU.EQ.0.0) GNU=1,E=10	MERGF149
	FISA(KK1)=XNUSIG(KK1)/GNU	MERGF150
	CAPA(KK1)=SABS(KK1)-FISA(KK1)	MERGF151
	PI TOT=0.0	MERGF152

3050	FORMAT (1P,E15,5)	MERGF153
	DO 3060 MGM=1,MGPS	MERGF154
	PITOT=PITOT+P*(MG,MGM)	MERGF155
3060	CONTINUE	MERGF156
	PRINT 3050,PITOT	MERGF157
	P0(KK1,KK1)=P0(KK1,KK1)-PITOT	MERGF158
	STR(KK1)=STR(KK1)-PITOT	MERGF159
	PRINT 3050,P0(KK1,KK1),STR(KK1)	MERGF160
3020	CONTINUE	MERGF161
	GO TO 3500	MERGF162
3000	CONTINUE	MERGF163
	PRINT 4115,NUCID(M),FID,(TTL(KK),KK=1,11)	MERGF164
1	FORMAT (A10)	MERGF165
	DO 10 K=1,NOBG	MERGF166
	KF=K-1	MERGF167
	IF (K.EQ.1) KF=NOBG2	MERGF168
	READ (NTPF,3) FISA(K),SABS(K),XNUSIG(K),STR(K),P0(K,K),P0(KF,K)	MERGF169
	READ (NTPF,3) ZILCH	MERGF170
	IF (ZILCH.FQ.0.0) GO TO 20	MERGF171
	PRINT 19,ZILCH	MERGF172
19	FORMAT (1H1,* ZILCH = *1PE12.5,* READ ERROR*)	MERGF173
20	CONTINUE	MERGF174
3	FORMAT (6E12.5)	MERGF175
	CAPA(K)=SABS(K)-FISA(K)	MERGF176
	IF (FISA(K).GT.0.0) AVNU(K)=XNUSIG(K)/FISA(K)	MERGF177
10	CONTINUE	MERGF178
3500	CONTINUE	MERGF179
500	CONTINUE	MERGF180
	LOD=-1	MERGF181
	DO 95 K=1,NOBG	MERGF182
	TOTP0(K)=0.	MERGF183
	TOTP1(K)=0.	MERGF184
	TOTRA(K)=0.	MERGF185
	TOTN2N(K)=0.	MERGF186
	SS(K)=0.	MERGF187
	XNUSIG(K)=FISA(K)*AVNU(K)	MERGF188
	DO 80 KF=K,NOBG2	MERGF189
	TOTRA(K)=TOTRA(K)+SINTRA(K,KF)	MERGF190
	IF (LOD.FQ.-1) SK2KF(K,KF)=P0(K,KF)+SINTRA(K,KF)+SSN2N(K,KF)*2.	MERGF191
	SS(K)=SS(K)+SK2KF(K,KF)	MERGF192
	TOTN2N(K)=TOTN2N(K)+SSN2N(K,KF)	MERGF193
80	CONTINUE	MERGF194
	CAPA(K)=CAPA(K)+SNP(K)+SNA(K)	MERGF195
	SABS(K)=CAPA(K)+FISA(K)	MERGF196
95	CONTINUE	MERGF197
	LOD1=LOD+2	MERGF198
	HOL(2)=10H ETOG=ENDF	MERGF199
	GO TO (110,120,130,140,150),LOD1	MERGF200
110	HOL(1)=10H ISOTROP T	MERGF201
	HOL(2)=10H RNSP TABLE	MERGF202
	GO TO 160	MERGF203
120	HOL(1)=10H P=0 TABLE	MERGF204
	GO TO 160	MERGF205
130	HOL(1)=10H P=1 TABLE	MERGF206
	GO TO 160	MERGF207
140	HOL(1)=10H P=2 TABLE	MERGF208
	GO TO 160	MERGF209
150	HOL(1)=10H P=3 TABLE	MERGF210
160	CONTINUE	MERGF211
	LOD=LOD+1	MERGF212
	L=LENG	MERGF213
	N=LENG-1	MERGF214

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C DOWN SCATTERING IS CONSIDERED ONLY BETWEEN ADJACENT GROUPS EXCEPT FOR MERGF215
C FIRST L GROUPS IN WHICH ALL CASES ARE CONSIDERED. ADDITIONAL VALUES AMERGF216
C ADDED INTO THE L -TH GROUP. MERGF217
DO 90 K=1,NOBG MERGF218
MN=LENG*K MERGF219
IF (MN.GT.NOBG*1) GO TO 90 MERGF220
DO 91 KF=MN,NOBG2 MERGF221
91 SK2KF(K,MN-1)=SK2KF(K,MN-1)+SK2KF(K,KF) MERGF222
90 CONTINUE MERGF223
MK=NOBG-LENG+2 MERGF224
DO 51 K1=MK,NOBG MERGF225
51 SK2KF(K1,NOBG)=SK2KF(K1,NOBG)+SK2KF(K1,NOBG2) MERGF226
300 CONTINUE MERGF227
LTABL=LENG+LNGUP+7 MERGF228
NXC=LTABL*NOBG MERGF229
NCX6=NXC+LTABL MERGF230
DO 360 JT=1,NCX6 MERGF231
360 A(JT)=0. MERGF232
DO 361 K=1,NOBG MERGF233
J1=(K-1)*LTABL+1 MERGF234
J2=J1+1 MERGF235
J3=J1+2 MERGF236
J4=J3+1 MERGF237
J5=J4+1 MERGF238
J6=J5+1 MERGF239
J7=J6+1 MERGF240
A(J1)=FISA(K) MERGF241
A(J4)=CAPA(K) MERGF242
A(J5)=SARS(K) MERGF243
A(J6)=XNUSIG(K) MERGF244
A(J7)=STR(K) MERGF245
DO 362 KF=K,NOBG2 MERGF246
J8=(KF-1)*LTABL+KF+7-K+1+LNGUP MERGF247
A(J8)=SK2KF(K,KF) MERGF248
IF ((KF-K).GE. LENG) GO TO 361 MERGF249
362 CONTINUE MERGF250
IF (FID.NE.6H 12.00) GO TO 361 MERGF251
IF (K.LT.KG1(NDKS)) GO TO 361 MERGF252
KUP1=K+1 MERGF253
IF (KUP1.GT.NOBG) GO TO 361 MERGF254
DO 3561 KF=KUP1,NOBG MERGF255
J9=LTABL*(K-1)+8+LNGUP-(KF-K) MERGF256
A(J9)=P0(KF,K) MERGF257
3561 CONTINUE MERGF258
361 CONTINUE MERGF259
250 CONTINUE MERGF260
PUNCH 85,NOBG,LTABL,MATNO,NUCID(M),(HOL(I),I=1,2) MERGF261
NP1=1 MERGF262
NBR=1 MERGF263
63 NP2=NP1+5 MERGF264
IF (NP2.LE.NXC) GOTO 67 MERGF265
PUNCH 65,(A(NP),NP=NP1,NP2),NUCID(M),NBR MERGF266
68 FORMAT (1P6E12.5) MERGF267
GO TO 64 MERGF268
67 PUNCH 65,(A(NP),NP=NP1,NP2),NUCID(M),NBR MERGF269
IF (NP2.EQ.NXC) GOTO 64 MERGF270
NBR=NBR+1 MERGF271
NP1=NP2+1 MERGF272
GO TO 63 MERGF273
64 CONTINUE MERGF274
65 FORMAT (1P6E12.5,A6,I2) MERGF275
DO 103 K=1,NOBG MERGF276

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103	KK(K)=K	MERGF277
	IF(M.NE.I)GOTO9	MERGF278
	WRITE(6,1000)(TITLE(I),I=1,7)	MERGF279
1000	FORMAT(8A10)	MERGF280
	9 CONTINUE	MERGF281
	WRITE(6,85) NOBG,LTABL,MATNO,NUCID(M),(HOL(I),I=1,2)	MERGF282
C	CHECK TO ADD UP SK2KF	MERGF283
C	THE DIAGONAL SUM OF THE DOWN SCATTERING AND SELF SCATTERING ADDED TO	MERGF284
C	ABSORPTION MUST EQUAL THE TRANSPORT CROSS SECTION.	MERGF285
	DO 400 K=1,NOBG	MERGF286
	DIAGSM(K)=0.	MERGF287
	KK=K+LENG-1	MERGF288
	IF(KK.GT.NOBG) KK=NOBG	MERGF289
	DO 450 KF=K,KK	MERGF290
450	DIAGSM(K)=DIAGSM(K)+SK2KF(K,KF)	MERGF291
	CHKSM(K)=DIAGSM(K)+SABS(K)-TOTN2N(K)	MERGF292
	DIFF(K)=CHKSM(K)-STR(K)	MERGF293
	WRITE(6,460)	MERGF294
460	FORMAT(20H0TRANSPORT 18HCHECKSUM	MERGF295
	1 10HDIFFERENCE)	MERGF296
	WRITE(6,470) STR(K),CHKSM(K),DIFF(K)	MERGF297
470	FORMAT(E12.4,7X,E12.4,8X,E12.4)	MERGF298
	IF(DIFF(K).LE.0.1*STR(K))GOTO400	MERGF299
	WRITE(6,480)	MERGF300
480	FORMAT(31HDIFFERENCE EXCEEDS 10 PER CENT)	MERGF301
400	CONTINUE	MERGF302
	WRITE(6,85) NOBG,LTABL,MATNO,NUCID(M),(HOL(I),I=1,2)	MERGF303
	DO 701 I=1,NOBG	MERGF304
	K=LTABL*(I-1)	MERGF305
	DO 701 J=1,LTABL	MERGF306
701	SK2KF(I,J)=A(K+J)	MERGF307
	KA=1 \$ KB=A	MERGF308
724	KC=MIN0(KB,NOBG)	MERGF309
	WRITE(6,720) (K,K=KA,KC)	MERGF310
	WRITE(6,721)	MERGF311
	DO 722 J=1,LTABL	MERGF312
722	WRITE(6,723) J,(SK2KF(I,J),I=KA,KC)	MERGF313
	KA=KA+8 \$ KB=KB+8	MERGF314
	IF(KA.LE.NOBG) GO TO 724	MERGF315
720	FORMAT(1H0,5X,8(2X,5HGROUP,13,3X))	MERGF316
721	FORMAT(1H)	MERGF317
723	FORMAT(14,1P8E13,5)	MERGF318
100	CONTINUE	MERGF319
	MATNO=MATNO+1	MERGF320
999	CONTINUE	MERGF321
	IF(HLT(1).NE.8HLASTDECK) GO TO 5002	MERGF322
	END	MERGF323

APPENDIX D

DANCOFF FACTOR FOR A REGULAR ARRAY OF CYLINDRICAL FUEL RODS

The Dancoff factor is an important quantity in the Levine method of space-shielding cross sections to account for the gross (fuel-rod) heterogeneity in the reactor core. For this purpose, a special computer program was written to calculate the Dancoff factor by three methods and to compare their results.

One method, due to Carlvik,²² gives the Dancoff factor by exact integration:

$$C = \frac{1}{\alpha_0} \int_0^{\alpha_0} d\alpha \frac{1}{2r} \int_{-r}^r dy \frac{Ki_3(t)}{Ki_e(0)} \quad , \quad (D-1)$$

where Ki_3 is the Bickley function, t is the optical length between rods, r is the radius of one rod, and α_0 is 30° for a hexagonal lattice.

For a hexagonal lattice, $y = r/d$, where d is the lattice pitch and y is related to the mean chord length $\bar{\ell}$ in the moderator through

$$\bar{\ell} \equiv \frac{4V_m}{S} = \frac{2d}{\pi y} \left(\frac{\sqrt{3}}{2} - \pi y^2 \right) \quad . \quad (D-2)$$

Sauer³⁰ has found a good approximation for C :

$$C = \frac{e^{-t\Sigma\bar{\ell}}}{1 + (1-t)\Sigma\bar{\ell}} \quad , \quad (D-3)$$

where Σ is the moderator cross section and, for a hexagonal lattice

$$t = \frac{\pi}{2} y \frac{1 - 2y}{\frac{\sqrt{3}}{2} - \pi y^2} - 0.12 \quad . \quad (D-4)$$

Bonalumi³¹ has pointed out that Sauer's Dancoff correction is bad approximation for very large moderator cross sections in the two cases of very large and very small volume ratios, i.e., for y near 0 and near 0.5.

Bonalumi has, therefore, suggested the following modification:

$$C = \frac{e^{-t\bar{\Sigma}\bar{\ell}}}{1 + (1 - t_1)\bar{\Sigma}\bar{\ell}} \quad , \quad (D-5)$$

where

$$t_1 = t + \frac{\bar{\Sigma}\bar{\ell}}{7 + \beta\bar{\Sigma}\bar{\ell}} \quad , \quad (D-6)$$

and $\beta = 2.125$ for a hexagonal lattice.

For the HTGR core configurations under study, all three methods of calculating the Dancoff factor have been found good, yielding very close answers. The listing of the computer program used for this comparison is included at the end of this appendix.

LASL Identification No. LP-0758

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PROGRAM DANCPIN(INP,OUT,PUN,FILM)                                DANCP  1
C  CALCULATES DANCOFF FACTOR FOR A REGULAR ARRAY OF INFINITE CYLINDERSDANCP  2
C  INPUT QUANTITIES                                             DANCP  3
C  NALF AND NRAD DETERMINE AN INTEGRATION GRID FOR THE CARLVIK  DANCP  4
C  INTEGRATION; THEY ARE BOTH TAKEN TO BE 128 .                DANCP  5
C  NLAT DETERMINES THE TYPE OF LATTICE , IT IS 4 FOR A SQUARE DANCP  6
C  LATTICE AND 6 FOR A HEXAGONAL LATTICE .                     DANCP  7
C  IF IOPTC=0, ALL THREE METHODS ARE COMPARED . IF IOPTC=1,    DANCP  8
C  THE CARLVIK ROUTINE IS USED , IF IOPTC=2, THE BONALUMI     DANCP  9
C  APPROXIMATION ONLY IS USED , IF IOPTC=3, THE SAUER         DANCP 10
C  APPROXIMATION ONLY IS USED .                                DANCP 11
C  RAD0 IS THE PIN RADIUS IN CM .                               DANCP 12
C  RAD1 IS THE MODERATOR RADIUS IN THE THREE-REGION MODEL .  DANCP 13
C  GAPWID IS THE GAP WIDTH (CM) AROUND THE PIN .             DANCP 14
C  RADIS IS THE MODERATOR RADIUS WHEN THE GAP-WIDTH IS NOT    DANCP 15
C  EXPLICITLY GIVEN AS IN THE SAUER OR THE BONALUMI APPROXIMATIONS . DANCP 16
C  SIGF IS THE MACROSCOPIC FUEL-PIN CROSS SECTION (1/CM) .   DANCP 17
C  DENSF IS THE ATOMIC CONCENTRATION OF THE FUEL PIN .       DANCP 18
C  SIGMAM IS THE MODERATOR MACROSCOPIC CROSS SECTION (1/CM) . DANCP 19
C  READ 5,NLAT,NALF,NRAD,IOPTC,RADIS                             DANCP 20
5  FORMAT(4I10,2E10.4)                                          DANCP 21
  IF(NLAT.NE.4.OR.NLAT.NE.6) PRINT 6                            DANCP 22
6  FORMAT(1H0,*NLAT MUST BE EQUAL TO 4 OR TO 6 , TRY AGAIN*)  DANCP 23
  IF(NRAD.LE.128) NRAD=128                                       DANCP 24
  IF(NALF.LE.128) NALF=128                                       DANCP 25
  READ 10,RAD0,RAD1,GAPWID,SIGMAM,SIGF,DENSF                    DANCP 26
10  FORMAT(6F12.6)                                              DANCP 27
  PRINT 15,NLAT,NALF,NRAD,IOPTC,RADIS                             DANCP 28
15  FORMAT(1H1,*NLAT = *,I4,* NALF = *,I4,* NRAD = *,I4,* IOPTC = *,DANCP 29
  1I4,* RADIS = *,E10.4)                                          DANCP 30
  PRINT 16,RAD0,RAD1,GAPWID,SIGMAM,SIGF,DENSF                    DANCP 31
16  FORMAT(1H0,*RAD0 = *,E12.6,* RAD1 = *,E12.4,* GAPWID = *,E12.6// DANCP 32
  1* SIGMAM = *,E12.6,* SIGF = *,E12.6,* DENSF = *,E12.6)      DANCP 33
  IF(IOPTC)20,20,30                                             DANCP 34

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20	DO 25 IOPTC=1,3	DANCP 35
	CALL DANC OFF(NLAT,RAD0,RAD1,GAPWID,SIGMAM,NALF,NRAD,CC,CS,	DANCP 36
	1CSB,C,IOPTC,RAD1S)	DANCP 37
	CALL LEVINE(SIGF,DEN SF,C,RAD0,IOPTC)	DANCP 38
25	CONTINUE	DANCP 39
	GO TO 50	DANCP 40
30	CALL DANC OFF(NLAT,RAD0,RAD1,GAPWID,SIGMAM,NALF,NRAD,CC,CS,	DANCP 41
	1CSB,C,IOPTC,RAD1S)	DANCP 42
	CALL LEVINE(SIGF,DEN SF,C,RAD0,IOPTC)	DANCP 43
50	CONTINUE	DANCP 44
	END	DANCP 45

	SUBROUTINE LEVINE(SIGF,DEN SF,C,RAD0,IOPTC)	LEVIN 1
C	CALCULATES THE EFFECTIVE GEOMETRIC SHIELDING CROSS SECTION BY THE	LEVIN 2
C	LEVINE METHOD AND USING THE OTTER APPROXIMATION FOR THE LEVINE	LEVIN 3
C	FACTOR	LEVIN 4
	ELBARF=2.*RAD0	LEVIN 5
	TAUF=SIGF*ELBARF	LEVIN 6
	IF(TAUF,LE.0.) PRINT 20	LEVIN 7
20	FORMAT(1H0,*TAUF IS G. OR E. TO ZERO*)	LEVIN 8
	IF(TAUF,GE.2.) GO TO 30	LEVIN 9
	ALEVI=1.5013*0.14879*TAUF**0.5-0.17226*TAUF	LEVIN 10
	GO TO 40	LEVIN 11
30	ALEVI=1.+1./TAUF-1./TAUF**3.	LEVIN 12
40	CONTINUE	LEVIN 13
	PRINT 50,ALEVI	LEVIN 14
50	FORMAT(1H0,*THE LEVINE FACTOR EQUALS *,1PE12,4)	LEVIN 15
	QDEN=ELBARF*(1.+(ALEVI-1.)*C)	LEVIN 16
	QNUM=ALEVI*(1.-C)	LEVIN 17
	EFXSEC=QNUM/QDEN/DEN SF	LEVIN 18
	IF(IOPTC,EQ.1) ANAME=10H(CARLVIK)	LEVIN 19
	IF(IOPTC,EQ.2) ANAME=10H(BONALUMI)	LEVIN 20
	IF(IOPTC,EQ.3) ANAME=10H(SAUER)	LEVIN 21
	PRINT 10,ANAME,EFXSEC	LEVIN 22
10	FORMAT(1H0,*EFFECTIVE SIGMA WITH DANC OFF FACTOR *,A10,* = *,	LEVIN 23
	11PE12,4)	LEVIN 24
	RETURN	LEVIN 25
	END	LEVIN 26

	SUBROUTINE DANC OFF(NLAT,RAD0,RAD1,GAPWID,SIGMAM,NALF,NRAD,CC,CS,	DANCO 1
1	CSB,C,IOPTC,RAD1S)	DANCO 2
C	CALCULATES DANC OFF FACTORS BY THE ORIGINAL METHOD AS USED BY	DANCO 3
C	CARLVIK, BY THE SAUER APPROXIMATION AND BY THE BONALUMI	DANCO 4
C	APPROXIMATION AND COMPARES THE RESULTS OF THE THREE	DANCO 5
	IF(IOPTC)10,10,20	DANCO 6
10	PRINT 15	DANCO 7
15	FORMAT(1H0,*IOPTC IS ZERO OR NEGATIVE,TRY AGAIN*)	DANCO 8
	RETURN	DANCO 9

20	IF (IOPTC-1) 10,30,40	DANCO 10
30	CALL CARLVJK(NLAT,RAD0,RAD1,GAPWID,SIGMAM,NALF,NRAD,CC)	DANCO 11
	PRINT 35,CC	DANCO 12
35	FORMAT(1H0,*DANCOFF FACTOR (CARLVJK) = *,E12.6)	DANCO 13
	C=CC	DANCO 14
	RETURN	DANCO 15
40	IF (IOPTC-2) 30,50,60	DANCO 16
50	CALL BONAL(RAD0,RAD1S,NLAT,SIGMAM,CSB)	DANCO 17
	PRINT 45,CSB	DANCO 18
45	FORMAT(1H0,*DANCOFF FACTOR (BONALUMI) = *,E12.6)	DANCO 19
	C=CSB	DANCO 20
	RETURN	DANCO 21
60	IF (IOPTC-3) 50,70,80	DANCO 22
70	CALL SAUER(RAD0,RAD1S,NLAT,SIGMAM,CS)	DANCO 23
	PRINT 55,CS	DANCO 24
55	FORMAT(1H0,*DANCOFF FACTOR (SAUER) = *,E12.6)	DANCO 25
	C=CS	DANCO 26
	RETURN	DANCO 27
80	PRINT 90	DANCO 28
90	FORMAT(1H0,*IOPTC IS GREATER THAN 3,TRY AGAIN*)	DANCO 29
	RETURN	DANCO 30
	END	DANCO 31

	SUBROUTINE CARLVJK(NLAT,RAD0,RAD1,GAPWID,SIGMAM,NALF,NRAD,CC)	CARLV 1
C	CALCULATES DANCOFF FACTORS BY THE ORIGINAL METHOD AS IMPLEMENTED	CARLV 2
C	BY CARLVJK	CARLV 3
	PI=3.141592654	CARLV 4
	GAM=0.0	CARLV 5
	IF (NLAT.F0.6) GAM=PI/6.0	CARLV 6
	CNST1=(1./PI)**0.5	CARLV 7
	CNST2=(3.**0.5/2.0)**0.5	CARLV 8
	IF (NLAT.F0.6) CNST1=CNST1*CNST2	CARLV 9
	PITCH=RAD1/CNST1	CARLV 10
	R=RAD0/PITCH	CARLV 11
	E=(RAD0+GAPWID)/PITCH	CARLV 12
	E2=E*E	CARLV 13
	CONST=2.0/(PI*NALF*NRAD)	CARLV 14
	SIG =SIGMAM*PITCH	CARLV 15
	CC=0.0	CARLV 16
	ISIG=10./SIG+1.	CARLV 17
	NROW=MIN0(100,ISIG)	CARLV 18
	I1=2*NRAD	CARLV 19
	DZ=R/NRAD	CARLV 20
	DALF=PI/(NLAT*NALF)	CARLV 21
	ALF=-0.5*DALF	CARLV 22
	DO 60 N=1,NALF	CARLV 23
	ALF=ALF+DALF	CARLV 24
	CAG=COS(ALF+GAM)	CARLV 25
	DX=COS(GAM)/CAG	CARLV 26
	DY=SIN(ALF)/CAG	CARLV 27
	T=SIN(ALF+GAM)/CAG	CARLV 28
	Z=-R-0.5*DZ	CARLV 29
	DO 50 I=1,I1	CARLV 30
	Z=Z+DZ	CARLV 31
	X=Z*T-SQRT(E2-Z*Z)	CARLV 32
	F=CAG-Z	CARLV 33
	IF (F.GE.F) GO TO 10	CARLV 34
	IF (F.LE.R) GO TO 40	CARLV 35
	X=X-2.0*SQRT(E2-F*F)	CARLV 36

10	Y=Z/CAG+1.0	CARLV 37
	DO 30 J=1,NROW	CARLV 38
	IY=Y+DY	CARLV 39
	Y=Y+DY-IY	CARLV 40
	X=X+DX	CARLV 41
	F=-CAG*Y	CARLV 42
	IF(F.LE.(-E))GO TO 12	CARLV 43
	IF(F.GE.(-R))GO TO 40	CARLV 44
	X=X-2.0*SQRT(E2-F*F)	CARLV 45
12	F=CAG+F	CARLV 46
	IF(F.GE.E)GO TO 30	CARLV 47
	IF(F.LE.R)GO TO 40	CARLV 48
	X=X-2.0*SQRT(E2-F*F)	CARLV 49
30	CONTINUE	CARLV 50
	GO TO 50	CARLV 51
40	X=X+F*T-SQRT(E2-F*F)	CARLV 52
	Q=SIG*X	CARLV 53
	CALL BKLY(Q,BIC3)	CARLV 54
	CC=CC+BIC3	CARLV 55
50	CONTINUE	CARLV 56
60	CONTINUE	CARLV 57
	CC=CONST*CC	CARLV 58
	RETURN	CARLV 59
	END	CARLV 60

	SUBROUTINE SAUER(RAD0,RAD1,NLAT,SIGMAM,CS)	SAUER 1
C	CALCULATES DANCOFF FACTORS BY THE SAUER APPROXIMATION	SAUER 2
	PI=3.141592654	SAUER 3
	RADRA=RAD1/RAD0	SAUER 4
	VOLRA=RADRA*RADRA-1.	SAUER 5
	VOLSQR=(1.+VOLRA)**0.5	SAUER 6
	IF(NLAT.FQ.4)TAU=((PI/4.)**0.5*VOLSQR-1.)/VOLRA-0.08	SAUER 7
	IF(NLAT.FQ.6)TAU=((PI/(3.**0.5*2.)**0.5*VOLSQR-1.)/VOLRA-0.12	SAUER 8
	IF(TAU)10,10,20	SAUER 9
10	PRINT 15	SAUER 10
15	FORMAT(1H0,*TAU IS ZERO,NLAT IS WRONG*)	SAUER 11
	RETURN	SAUER 12
20	ELBARF=2.*RAD0	SAUER 13
	ELBARM=ELBARF*VOLRA	SAUER 14
	PROD=SIGMAM*ELBARM	SAUER 15
	DANCOF=EXP(-TAU*PROD)/(1.+(1.-TAU)*PROD)	SAUER 16
	CS=DANCOF	SAUER 17
	RETURN	SAUER 18
	END	SAUER 19

<pre> SUBROUTINE BONAL(RAD0,RAD1,NLAT,SIGMAM,CSB) C CALCULATES DANCOFF FACTORS BY THE BONALUMI APPROXIMATION TAU=0. PI=3.141592654 RADRA=RAD1/RAD0 VOLRA=RADRA*RADRA-1. VOLSQR=(1.+VOLRA)**0.5 IF(NLAT.EQ.4)TAU=((PI/4.)**0.5*VOLSQR-1.)/VOLRA IF(NLAT.EQ.6)TAU=((PI/(3.**0.5*2.))**0.5*VOLSQR-1.)/VOLRA IF(TAU)10,10,20 10 PRINT 15 15 FORMAT(1H0.*TAU IS ZERO,NLAT VALUE IS WRONG*) RETURN 20 ELBARF=2.*RAD0 ELBARM=ELBARF*VOLRA PROD=SIGMAM*ELBARM IF(NLAT.EQ.4)BETA=5.67 IF(NLAT.EQ.6)BETA=2.125 DELTAU=PROD/(7.+BETA*PROD) TAU1=TAU+DELTAU DANCOF=EXP(-TAU*PROD)/(1.+(1.-TAU)*PROD) CSB=DANCOF RETURN END </pre>	<pre> BONAL 1 BONAL 2 BONAL 3 BONAL 4 BONAL 5 BONAL 6 BONAL 7 BONAL 8 BONAL 9 BONAL 10 BONAL 11 BONAL 12 BONAL 13 BONAL 14 BONAL 15 BONAL 16 BONAL 17 BONAL 18 BONAL 19 BONAL 20 BONAL 21 BONAL 22 BONAL 23 BONAL 24 </pre>
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<pre> SUBROUTINE BKLY(X,BIC3) C CALCULATES BICKLEY FUNCTIONS OF THE THIRD ORDER A0=0.9379388841 A1=1.194191634 A2=0.588245154 A3=0.570337193 A4=-1.5791166 A5=4.292469 B0=0.7276787064 B1=0.9254690857 B2=0.4741520763 B3=0.250820355 B4=-0.025930075 B5=0.055707999 C0=0.4166740874 C1=0.5295655111 C2=0.2754273045 C3=0.1283775092 C4=0.0119191487 C5=0.0139209543 D0=0.2215940159 D1=-0.09388379097 D2=0.0147382145 D3=-0.000857650032 E0=0.2826723681 E1=0.2356320335 F2=0.06340205186 E3=0.01360032364 </pre>	<pre> BKLY 1 BKLY 2 BKLY 3 BKLY 4 BKLY 5 BKLY 6 BKLY 7 BKLY 8 BKLY 9 BKLY 10 BKLY 11 BKLY 12 BKLY 13 BKLY 14 BKLY 15 BKLY 16 BKLY 17 BKLY 18 BKLY 19 BKLY 20 BKLY 21 BKLY 22 BKLY 23 BKLY 24 BKLY 25 BKLY 26 BKLY 27 BKLY 28 </pre>
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F0=1.012074180	BKLY 29
F1=-0.000325432	BKLY 30
F2=-1.1646323	BKLY 31
F3=1.3873864	BKLY 32
F4=-4.4655208	BKLY 33
X2=X*X	BKLY 34
X3=X2*X	BKLY 35
X4=X3*X	BKLY 36
X5=X4*X	BKLY 37
IF(X)10,20,20	BKLY 38
10 PRINT 15	BKLY 39
15 FORMAT(1H0.*X IS LESS THAN ZERO,TRY AGAIN*)	BKLY 40
RETURN	BKLY 41
20 IF(X-0.1)30,40,40	BKLY 42
30 SUM=A0+A1*X+A2*X2+A3*X3+A4*X4+A5*X5	BKLY 43
BIC3=0.7366554521/SUM	BKLY 44
RETURN	BKLY 45
40 IF(X-0.4)50,60,60	BKLY 46
50 SUM=B0+B1*X+B2*X2+B3*X3+B4*X4+B5*X5	BKLY 47
BIC3=0.5714977571/SUM	BKLY 48
RETURN	BKLY 49
60 IF(X-1.0)70,80,80	BKLY 50
70 SUM=C0+C1*X+C2*X2+C3*X3+C4*X4+C5*X5	BKLY 51
BIC3=0.3272473766/SUM	BKLY 52
RETURN	BKLY 53
80 IF(X-2.5)90,100,100	BKLY 54
90 BIC3=(D0+D1*X+D2*X2+D3*X3)/(E0+E1*X+E2*X2+E3*X3)	BKLY 55
RETURN	BKLY 56
100 Y=1./(X+3.25)	BKLY 57
SUM=F0+F1*X+F2*X2+F3*X3+F4*X4	BKLY 58
BIC3=1.268445824*EXP(-X)/(Y**0.5*SUM)	BKLY 59
RETURN	BKLY 60
END	BKLY 61

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